

Addressing the impact of fishbones on core turbulence with gyrokinetic simulations

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INTRODUCTION AND MOTIVATIONS

In order to reach fusion-relevant temperatures in the core of fusion devices more easily, a reduction of outward fluxes would be beneficial. Since these fluxes are mainly driven by **microturbulence** (e.g. TEM, ITG) [1], the study of the interaction of these modes with **macroscopic plasma instabilities** (such as fishbones) is of fundamental importance.

Fishbones (FB) [2] develop around low order rational surfaces with $nq = m$, with q , n and m safety factor and toroidal and poloidal mode numbers. Recent results [3,4] show a link between FB trigger and turbulence level reduction. We investigate this physics with **GENE linear** and **nonlinear global simulations**.

TOOLS: THE GENE CODE

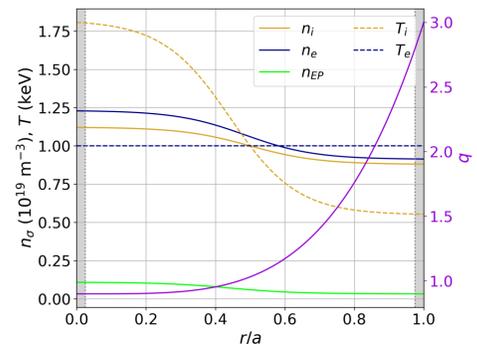
GENE [5] is a **eulerian gyrokinetic code** that solves the Vlasov-Maxwell system, composed by the **gyrokinetic equation**,

$$\frac{\partial F_\sigma}{\partial t} + \left[v_{\parallel} \hat{b}_0 + \frac{B_0}{B_{0\parallel}} (\vec{v}_{\vec{x}} + \vec{v}_{\nabla B} + \vec{v}_c) \right] \cdot \left[\vec{\nabla} F_\sigma - (q_\sigma \vec{\nabla} \phi_1 + \frac{q_\sigma}{c} \hat{b}_0 \dot{A}_{1\parallel} + \mu \vec{\nabla} B_0) \frac{1}{m_\sigma v_{\parallel}} \frac{\partial F_\sigma}{\partial v_{\parallel}} \right] = 0$$

along with the **equations for the evolution of the potentials** ϕ_1 and $A_{1\parallel}$ in the 5D phase space described by coordinates $\{x, y, z, v_{\parallel}, \mu\}$. We can perform **linear runs** by neglecting the nonlinear interaction term $\propto \vec{v}_{\vec{x}}$.

PLASMA PROFILES AND SETUP

Fishbones are triggered by the presence of **energetic particles** (EP) inside the tokamak. We consider a collisionless three species plasma: ions (H), electrons and EP (high- T D). Electron and EP temperatures are flat with $T_e = T_i(0.5) = 1$ keV and densities $n_i = 1$, $n_{EP} = 0.06$ and $n_e = 1.06 \times 10^{19} \text{m}^{-3}$ hold at $r/a = 0.5$. This setup is chosen in order to have only **$n=1$ FB** and **ITG** as unstable modes.



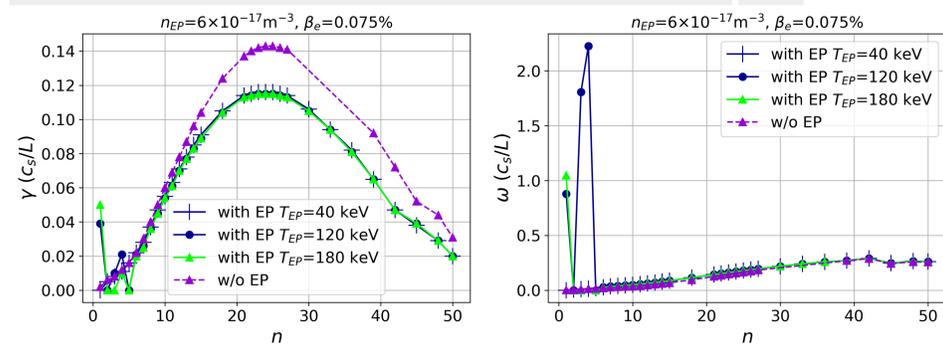
We consider a circular geometry with $R=10$ m and $a=1$ m. On-axis magnetic field is $B_T=1$ T, safety factor profile $q(r) = 0.9 + 2.1(r/a)^4$. **$q=1$ at $r/a=0.47$** holds. **$\beta_e = 0.075\%$** and the simulation domain is **$0.025 \leq r/a \leq 0.975$** .



LINEAR SIMULATIONS

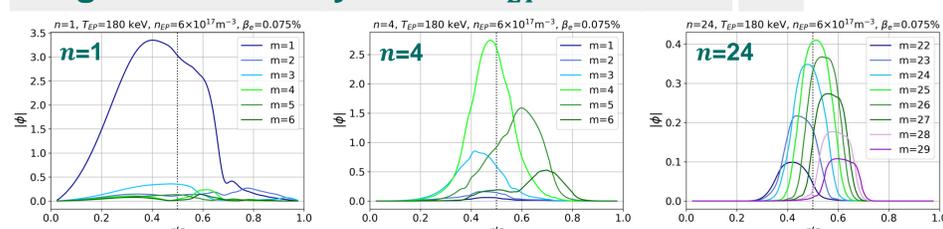
An analysis of the modes which develop in the setup is performed with GENE linear simulations. **Different flat EP temperatures** are considered.

Linear mode spectra



- An **ITG mode** branch peaks around $n=24$ for each setup.
- At **$T_{EP}=180$ keV** the $n=1$ mode is the only MHD instability.
- When **$T_{EP}=120$ keV**, $n=3$ and 4 are unstable along with the $n=1$ mode.
- The cases with **$T_{EP}=40$ keV** and **without EP** show no low- n mode. A **dilution effect** [6] of EP is observed in γ when removing them and $n_i \rightarrow n_e$.

Single mode analysis for $T_{EP}=180$ keV



- The **$m=1$** structure is clearly dominant for the $n=1$ mode.
- Different m values contribute to the mode for **higher n** .

NONLINEAR SIMULATIONS

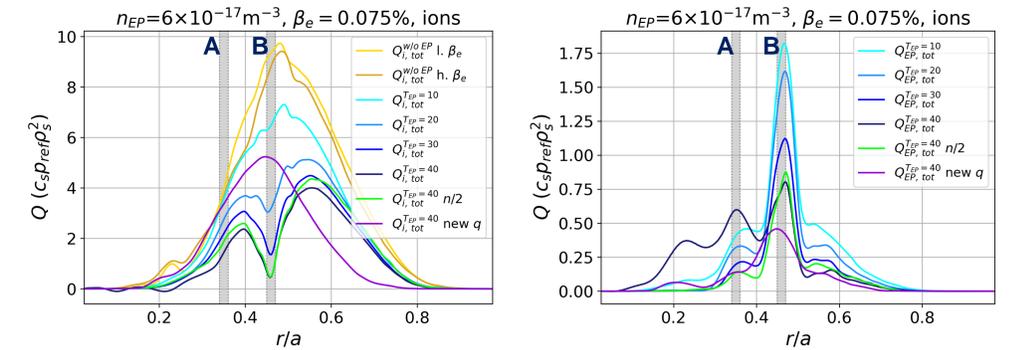
3

GENE nonlinear simulations are performed including **$n=0-47$ modes** for the cases **$T_{EP}=10-40$ keV**, and **w/o EP**, with same β_e ('l. β_e ') and β_{tot} ('h. β_e ') of the case with EP. A run with $T_{EP}=40$ and **$n=0,2,\dots,46$** is also considered, along with a setup with **shifted q profile**, ($r/a|_{q=1}=0.7$) do not coincide with the gradients maximum.

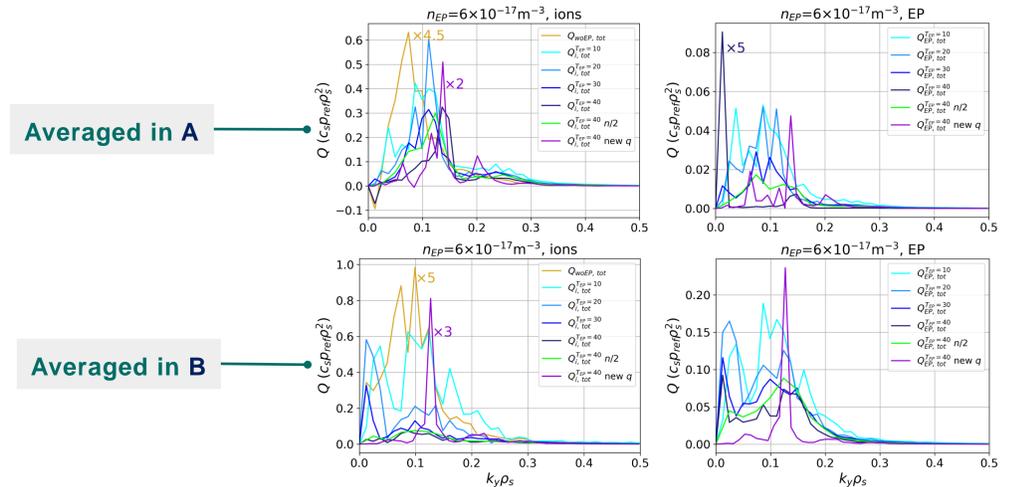
Flux radial profiles and spectra

3.1

Total (ES+EM) **heat fluxes Q_{tot}** are studied vs the radial position r/a and the wave number $k_y \rho_s$ (in regions **A** and **B**) for ions and EP.



- A reduction of fluxes is observed with increasing **T_{EP}** , in particular around **$q=1$** . No relevant effect of removing the **$n=1$ mode** can be observed in ion fluxes. Shifting the **q profile**, the ion flux depression in **B** disappears.

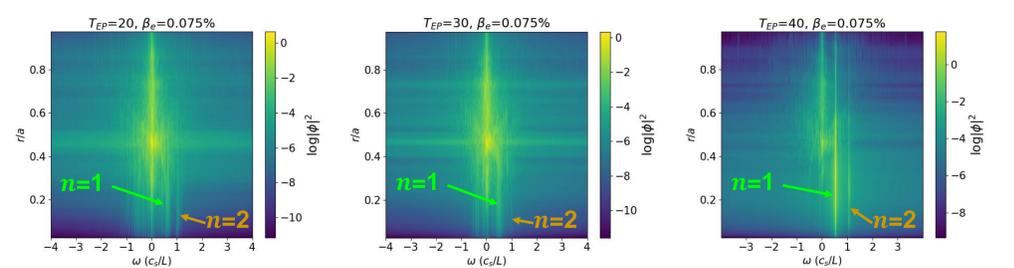


- In **A** the **$n=1$ mode** is the main driver of EP transport for **$T_{EP}=40$ keV** with ref. q .
- In **B**, **low- n modes** concur with the ITG branch to drive transport when EP are included and $q=1$ coincides with the maximum of gradients.

Frequency analysis

3.2

Spectrograms for **$n=1,2$** are computed vs r/a for **$T_{EP}=20, 30, 40$ keV**.



- The **$n=1$ mode intensifies** as **T_{EP} increases**, along with an expansion to higher r/a values of the radial domain of both modes.
- No definite frequencies can be identified for the **$n=1,2$ modes** in the setup with **$T_{EP}=40$ keV** and shifted q (spectrogram not reported here).

CONCLUSIONS AND PERSPECTIVES

- **Linearly** low- n high frequency modes are unstable for **high enough T_{EP}** values.
- **Nonlinearly**, ion fluxes are reduced by the inclusion of EP. The flux depression around the **$q=1$** position increases with the value of **T_{EP}** .
- **Low- n modes** are the main drivers of EP transport, with an **$n=1$ dominance** in **A** for **$T_{EP}=40$** . Their contribution is suppressed when the q profile is shifted.
- As **T_{EP} increases**, **$n=1,2$ modes** expand towards region **B** and the $n=1$ mode becomes more and more relevant. It is not clear yet if these modes are linked to the ion heat flux reduction observed at $q=1$ with increasing **T_{EP}** . Further analysis (e.g. bicoherence) and simulations (cases with **$T_{EP}>40$ keV**) are ongoing.
- **$n=1$ runs** with **ORB5** [7] are ongoing for comparison with GENE results.

[1] W Horton, 1999 *Rev. Mod. Phys.* **71** 3

[2] F Zonca et al, 2009 *Nucl. Fusion* **49** 085009

[3] X X He et al, 2022 *Plasma Phys. Control. Fusion* **64** 015007

[4] G Brochard et al, 2024 *Phys. Rev. Lett.* **132** 075101

[5] F Jenko et al, 2000 *Phys. Plasmas* **7** 1904

[6] A Di Siena et al, 2018 *Nucl. Fusion* **58** 054002

[7] E Lanti et al, 2020 *Comput. Phys. Commun.* **251** 107072

