



Dynamics of energetic particles in the vicinity of Black Holes

TOK/MHD Retreat 2025

Daniele Villa

The fundamental problem

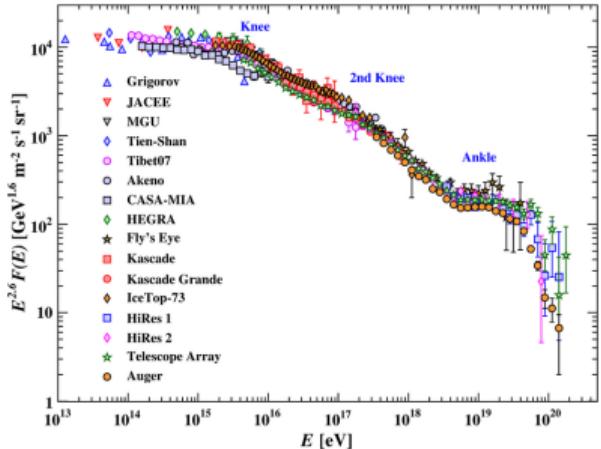


Figure: Source: CERN Document Server.

- Mechanism that **accelerates particles to** $\sim \text{EeV} = 10^{18} \text{eV}$ energies unknown \rightarrow must be of **extragalactic origin**.
- Such particles offer a **window into remote and extreme environments**.
- Black holes** “contain” huge amounts of energy \rightarrow if transferred to particles these could easily reach high energies.
- Understanding the **large-scale environment around black holes**.



Active Galactic Nuclei

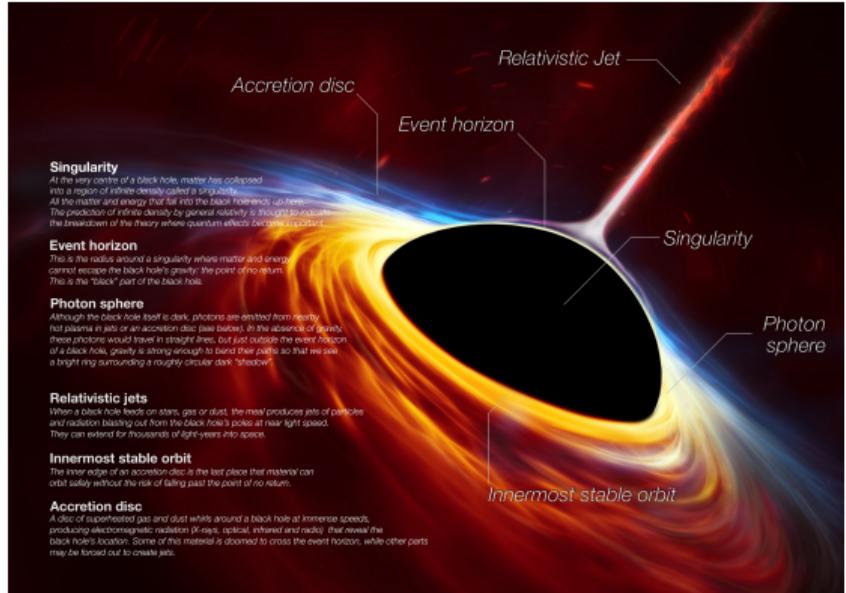


Figure: Source: ESO.

Active Galactic Nuclei

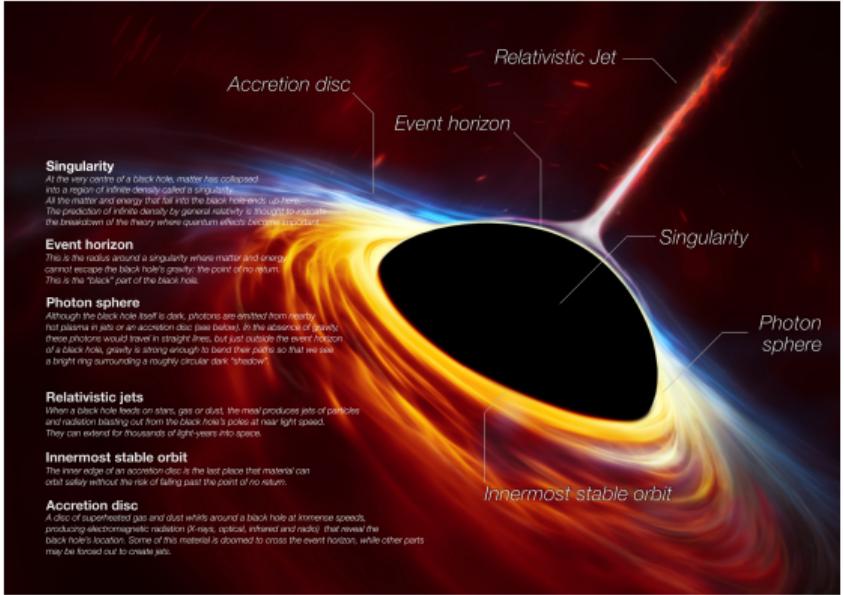


Figure: Source: ESO.

Central black hole →
Supermassive
($M \geq 10^6 M_{\odot}$), often
spinning

Active Galactic Nuclei

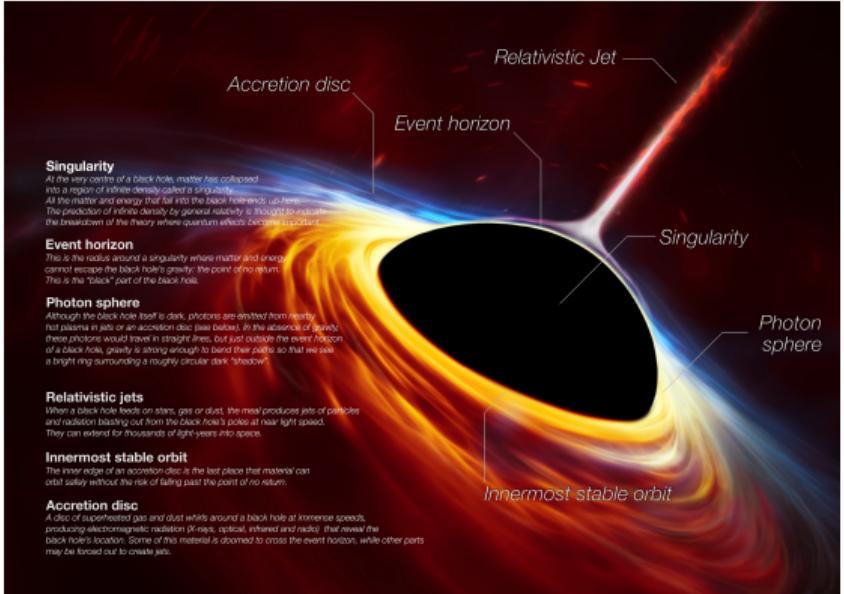


Figure: Source: ESO.

· **Central black hole** → Supermassive ($M \geq 10^6 M_{\odot}$), often spinning

· **Accretion Disk** → Matter captured by the gravitational well of the black hole → magnetized turbulent fluid

Active Galactic Nuclei

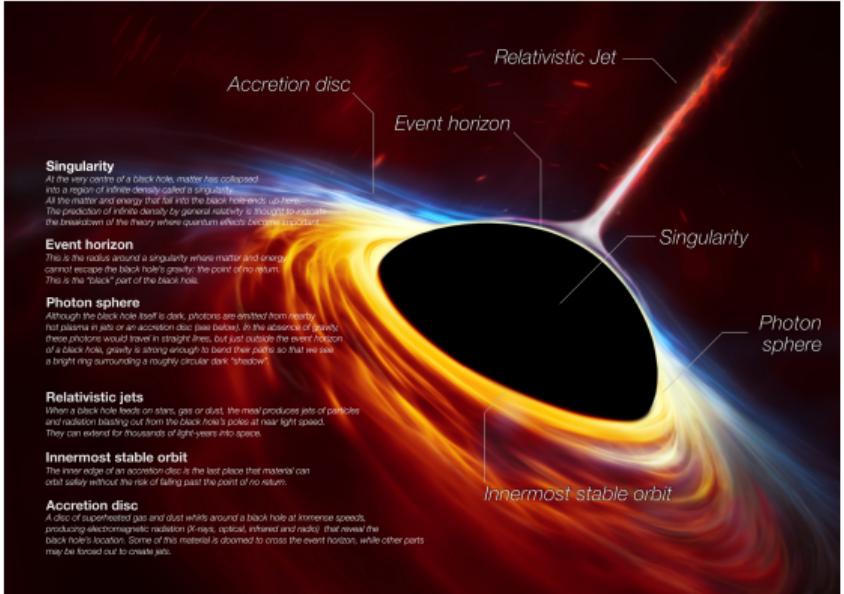


Figure: Source: ESO.

Central black hole → Supermassive ($M \geq 10^6 M_{\odot}$), often spinning

Accretion Disk → Matter captured by the gravitational well of the black hole → magnetized turbulent fluid

Relativistic Jet → Strongly accelerated (and strongly magnetized) outflow of matter

General Relativistic MHD

Ideal system of equations (*Y. Mizuno and L. Rezzolla, 2024, arXiv:2404.13824*)

$$\partial_t(\sqrt{\gamma}D) + \partial_i(\sqrt{\gamma}D(\alpha v^i - \beta^i)) = 0 \quad (1)$$

$$\partial_t(\sqrt{\gamma}S_j) + \partial_i(\sqrt{\gamma}(\alpha W_j^i - \beta^i S_j)) = \frac{1}{2}\sqrt{-g}T^{\mu\nu}\partial_j T_{\mu\nu} \quad (2)$$

$$\partial_t(\sqrt{\gamma}U) + \partial_i(\sqrt{\gamma}(\alpha S^i - U\beta^i)) = -\sqrt{-g}T^{\mu\nu}\nabla_\mu n_\nu \quad (3)$$

$$\partial_t(\sqrt{\gamma}B^j) + \partial_i(\sqrt{\gamma}B^j(\alpha v^i - \beta^i) - \sqrt{\gamma}B^i(\alpha v^j - \beta^j)) = 0 \quad (4)$$

Code: AthenaK (*J. M. Stone et al., 2024, arXiv:2409.16053*) (KOKKOS for parallel hybrid-architecture execution, i.e. multi-gpu support).

Adding particles in AthenaK

Development of a new physics module for AthenaK that allows single particle dynamics.
Currently only **tracer/pассив particles**.

Gravitational dynamics → iterative method for the free-particle Hamiltonian (*F. Bacchini et al., 2019, The Astrophysical Journal Supplement Series, 240, 2.*):

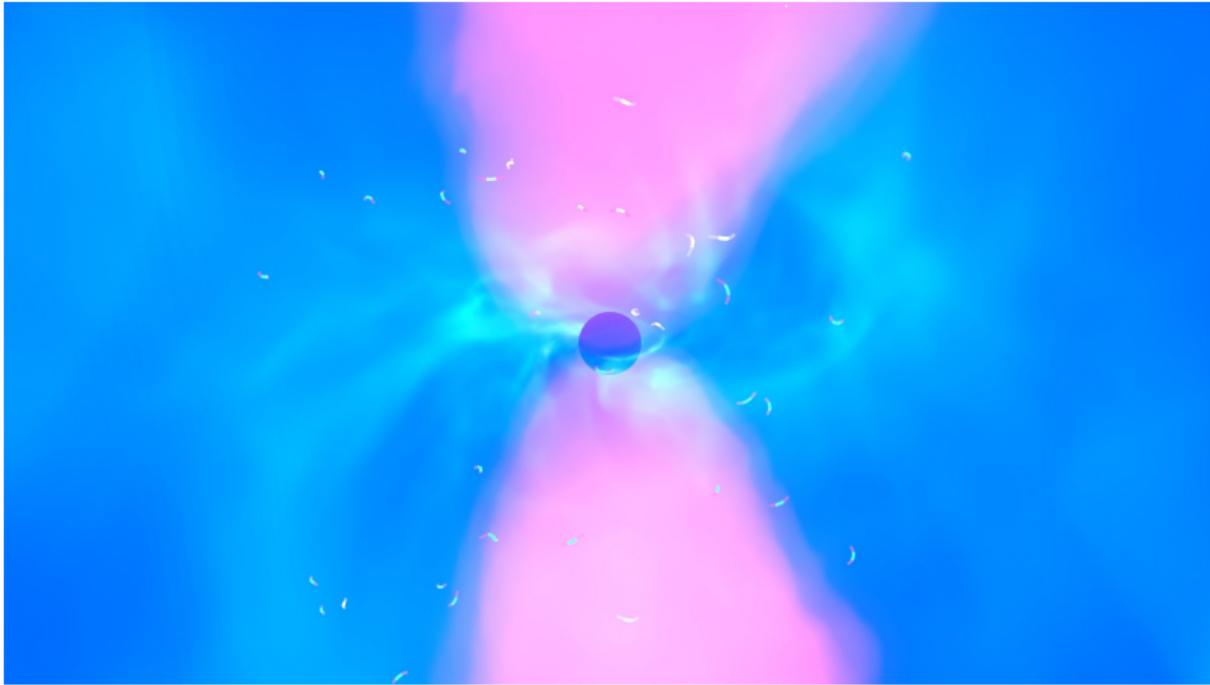
$$\mathcal{H}(x^i, u_i) = \alpha(\gamma^{jk}(x^i) u_j u_k + \epsilon)^{1/2} - \beta^j(x^i) u_j \quad (5)$$

Electromagnetic dynamics → explicit Boris steps adapted to GR (*B. Ripperda et al., 2018, The Astrophysical Journal Supplement Series, 235, 1.*)

$$\frac{u_i(t + \delta t) - u_i(t)}{\delta t} = \frac{q}{m} (E_i(x(t + \frac{1}{2}\delta t)) + \bar{v}_j \times B_k(x(t + \frac{1}{2}\delta t))) \quad \bar{v}_j = \frac{u_j(t + \delta t) + u_j(t)}{2\Gamma(t + \frac{1}{2}\delta t)} \quad (6)$$

Main hurdle at the moment is finding appropriate parameters for system of interest.

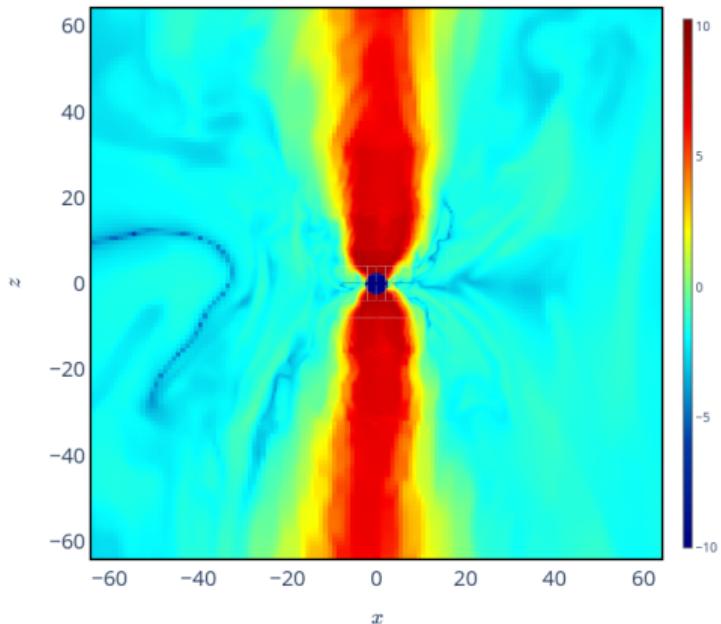
What does this “look like”?



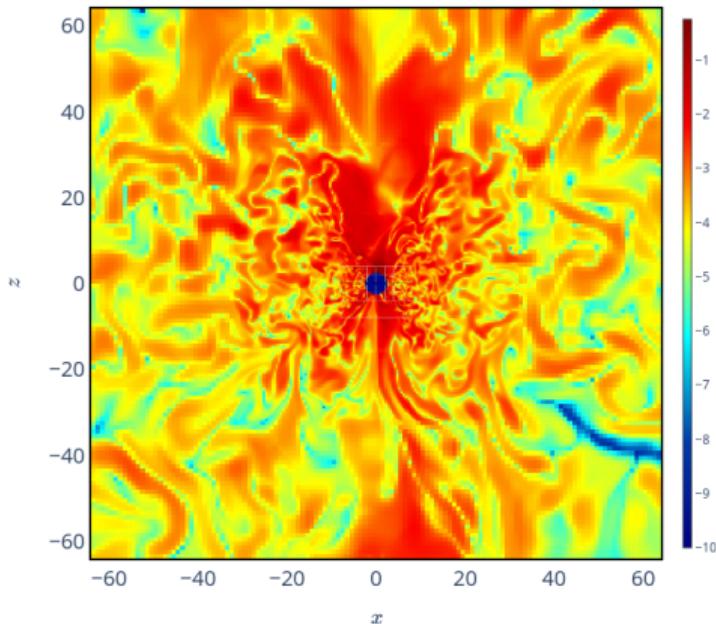
Accretion Disk Regimes

View IN the disk plane

MAD $a = 0.99$ $\log_{10}(\sigma)$



SANE $a = 0.99$ $\log_{10}(\sigma)$



Accretion Disk Regimes

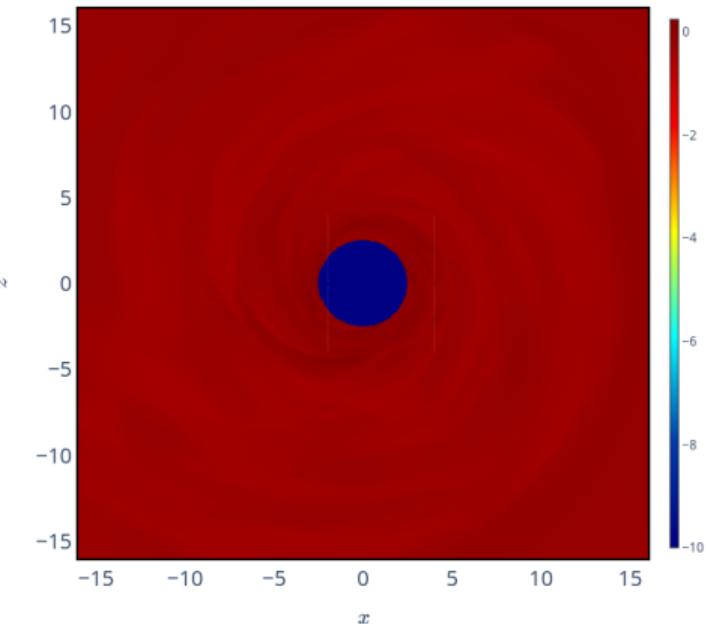
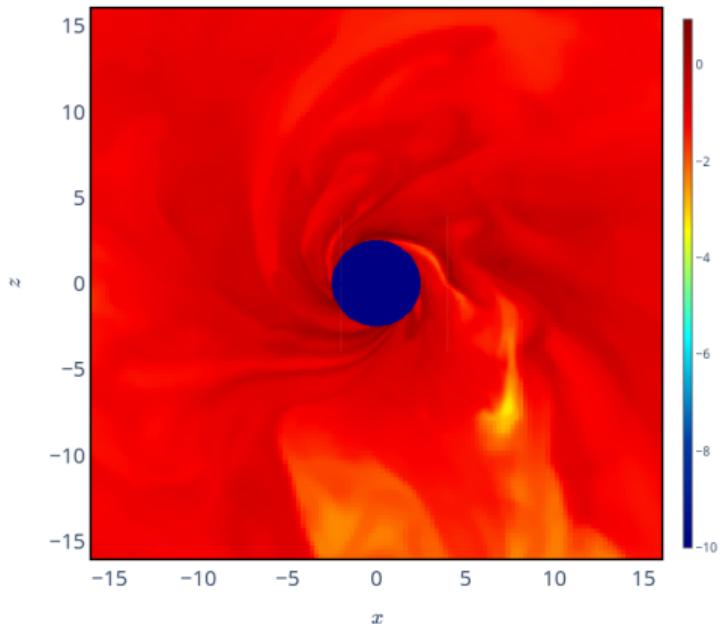
View **PERPENDICULAR TO** the disk plane

MAD $a = 0.99$

$\log_{10}(n)$

SANE $a = 0.99$

$\log_{10}(n)$

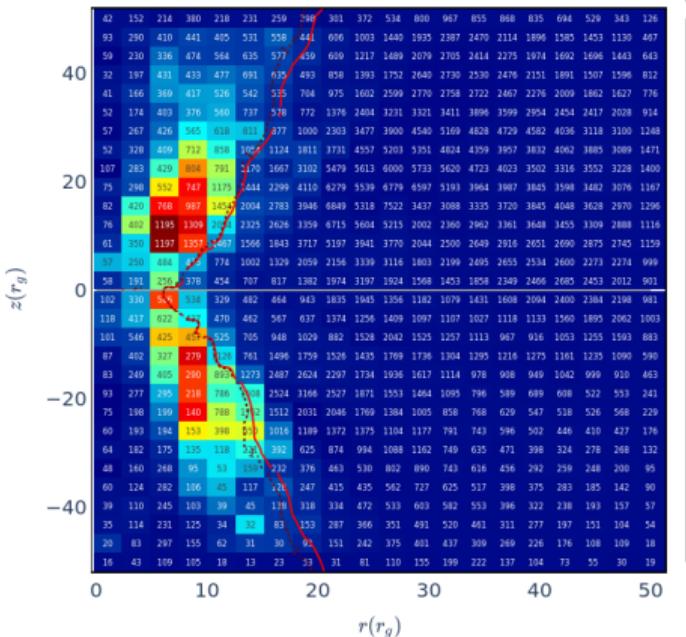




Average Particle Dynamics

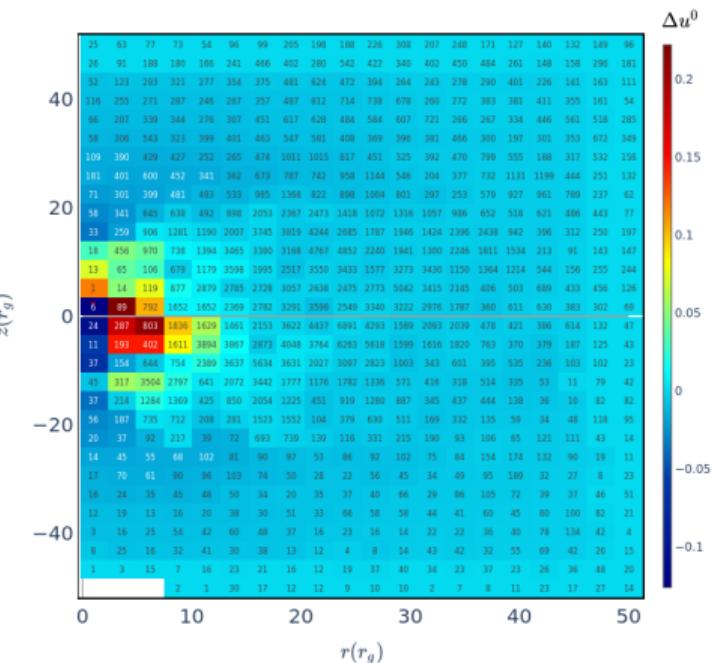
MAD $a = 0.99$

$\langle \Delta u^0 \rangle$



SANE $a = 0.99$

$\langle \Delta u^0 \rangle$

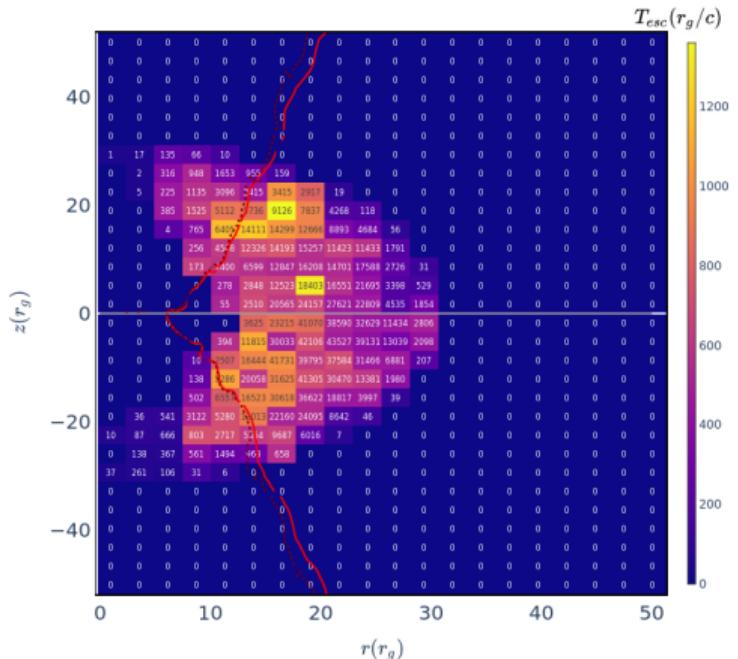




Average Particle Dynamics

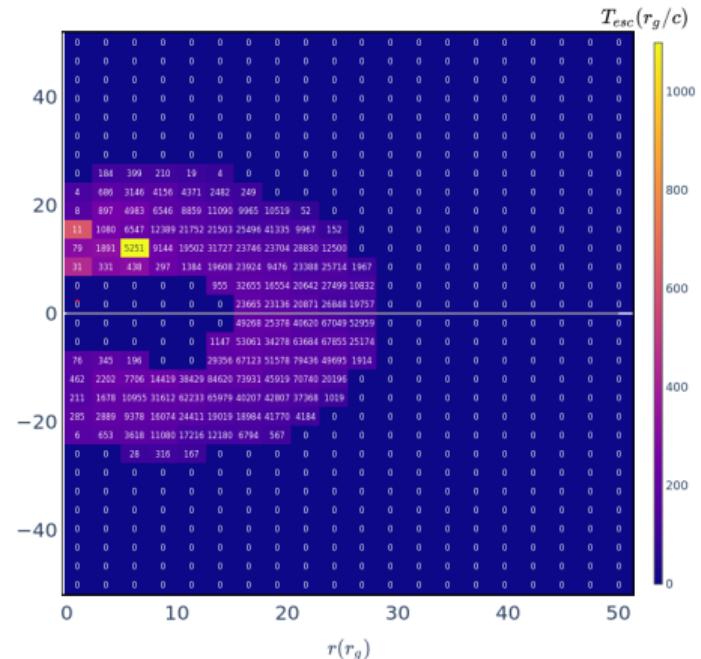
MAD $a = 0.99$

$$\langle \tau_{esc,out}^{max} \rangle$$



SANE $a = 0.99$

$$\langle \tau_{esc,out}^{max} \rangle$$

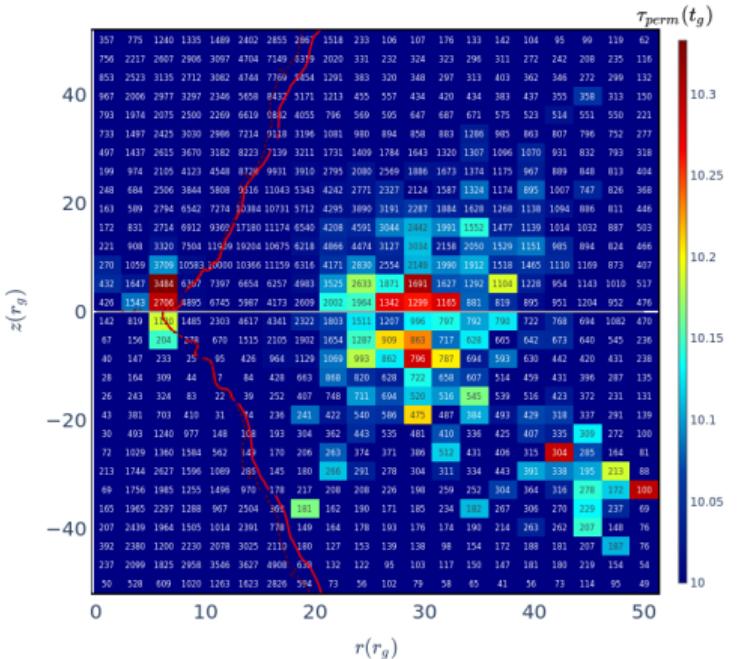




Average Particle Dynamics

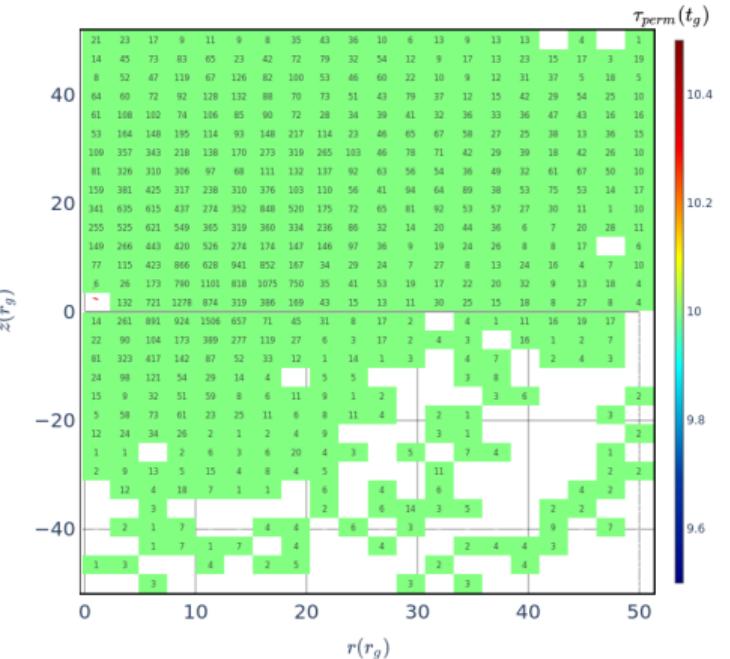
MAD $a = 0.99$

$\langle \tau_{trap} \rangle$



SANE $a = 0.99$

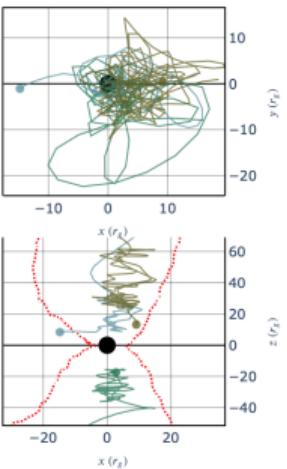
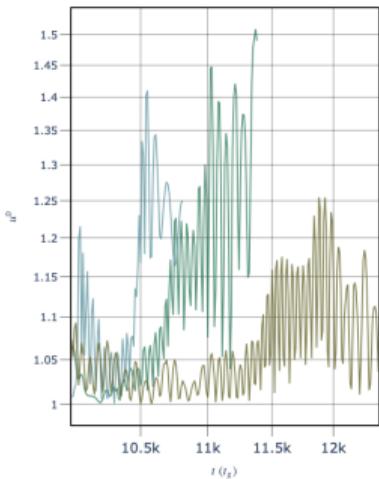
$\langle \tau_{trap} \rangle$



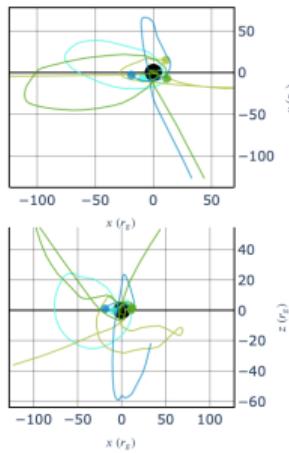
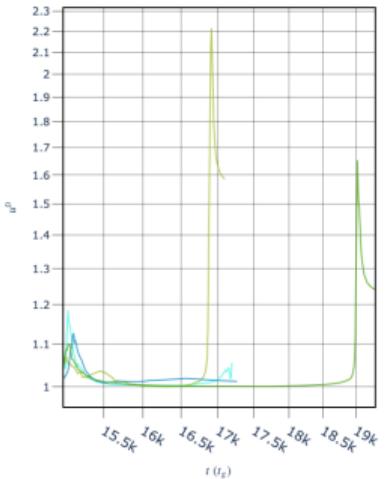


Individual Particle Dynamics

MAD $a = 0.99$



SANE $a = 0.99$





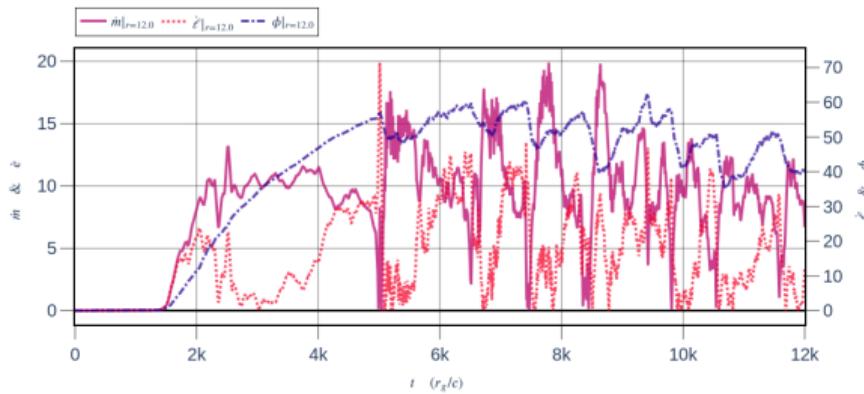
Outlook and future developments

- Use **particle pusher that evolves equations in Gyro-Center Approximation** to deal with system scale → Implemented, ongoing validation.
- Assess whether **current models of Accretion Disks are satisfactory** to explain neutrino emission from NGC 1068
 - Extend to **local resistivity** model.
 - Include **cooling and possibly radiation transport** in GRMHD simulations.
- Study the properties of **turbulence in GRMHD**.
- Study the **AGN NGC 1068 as a potential neutrino source**, trying to provide a first-principle explanation for the results from IceCube (*IceCube Collaboration, 2022, Science, 378, 6619.*)

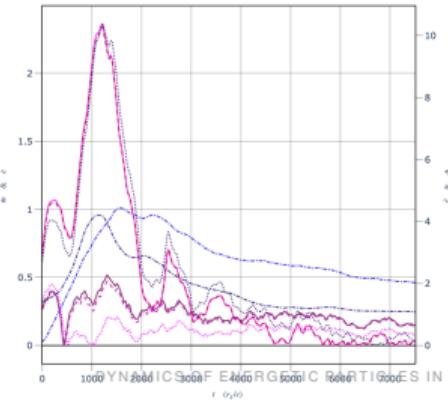
Extra slides

Accretion Disk Regimes

MAD regime



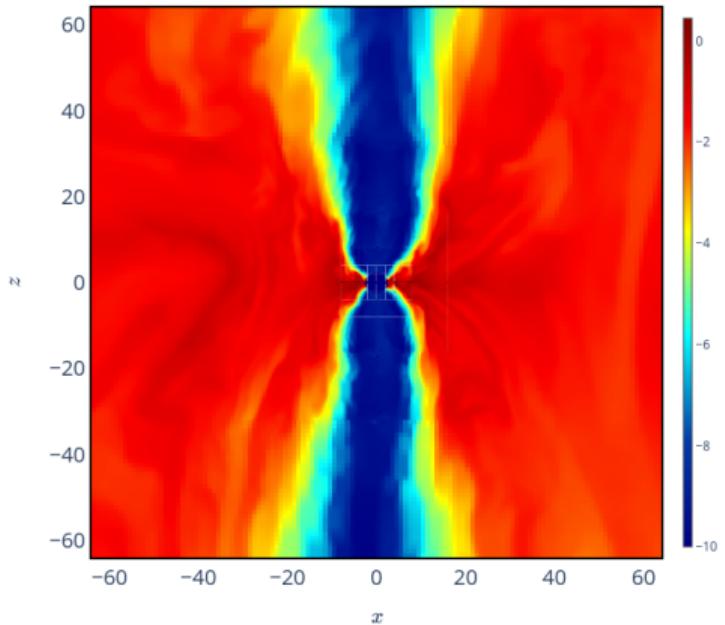
SANE regime



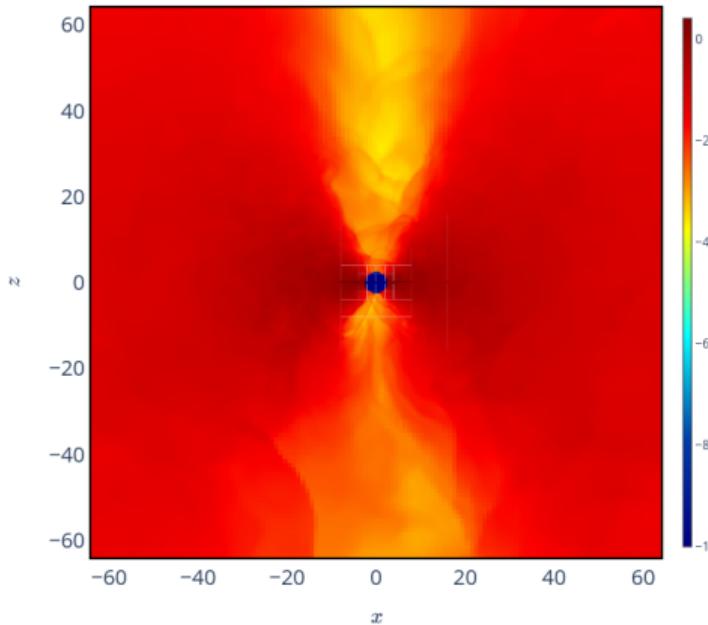
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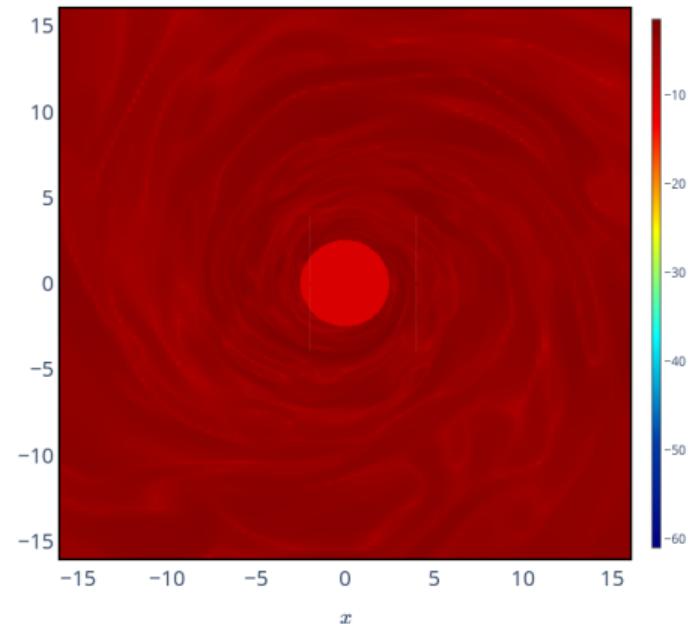
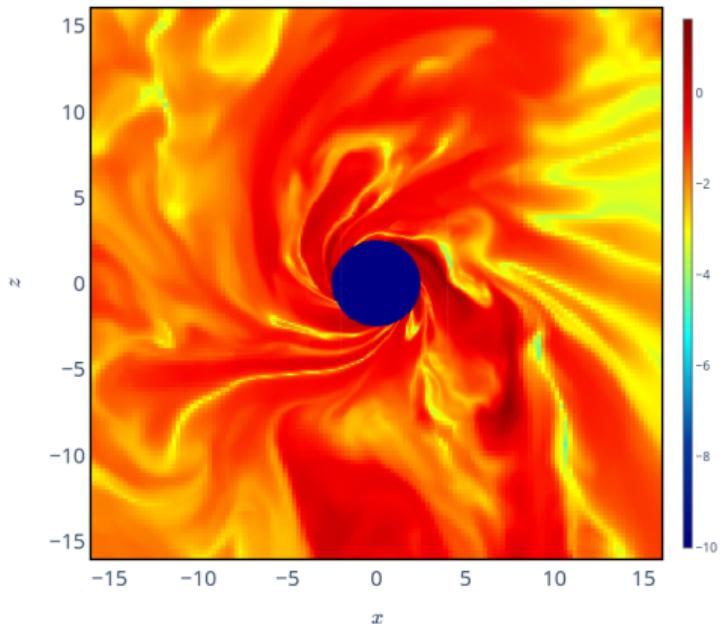
View **PERPENDICULAR TO** the disk plane

MAD $a = 0.99$

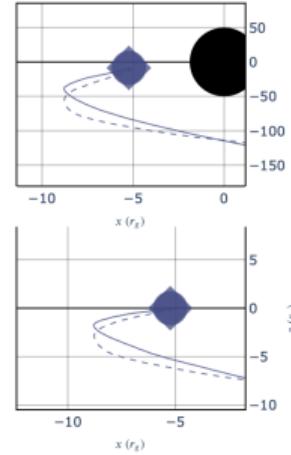
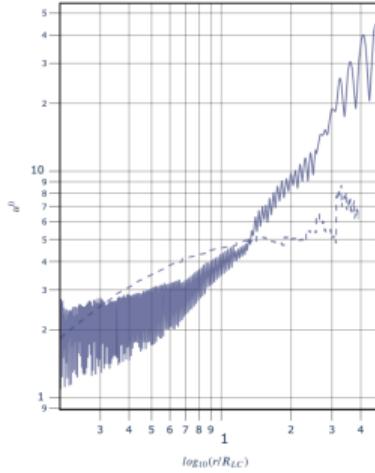
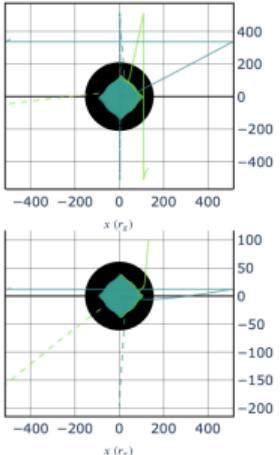
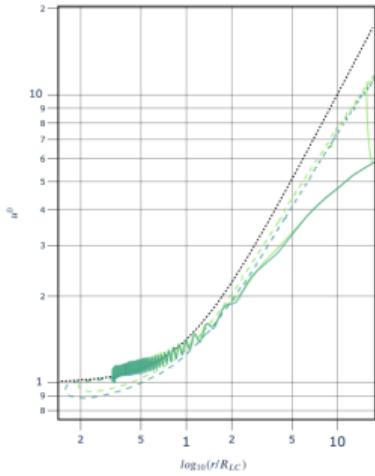
$\log_{10}(\sigma)$

SANE $a = 0.99$

$\log_{10}(\sigma)$



GCA vs Kinetic



Comparison of individual particle dynamics in MAD and SANE cases

Neutrino Measurements At IceCube

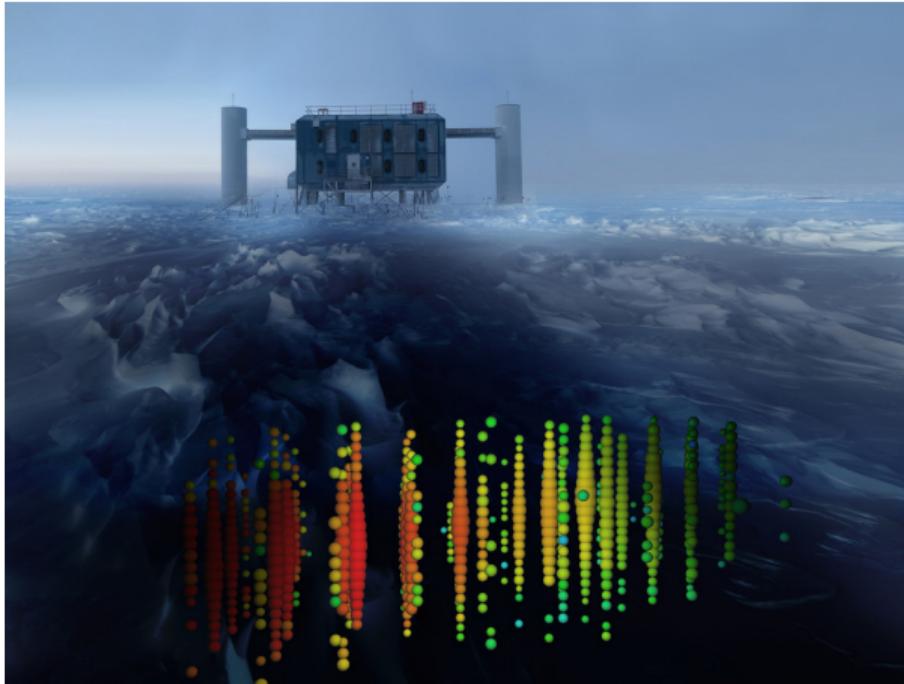


Figure: The IceCube Observatory.

Neutrino Measurements At IceCube

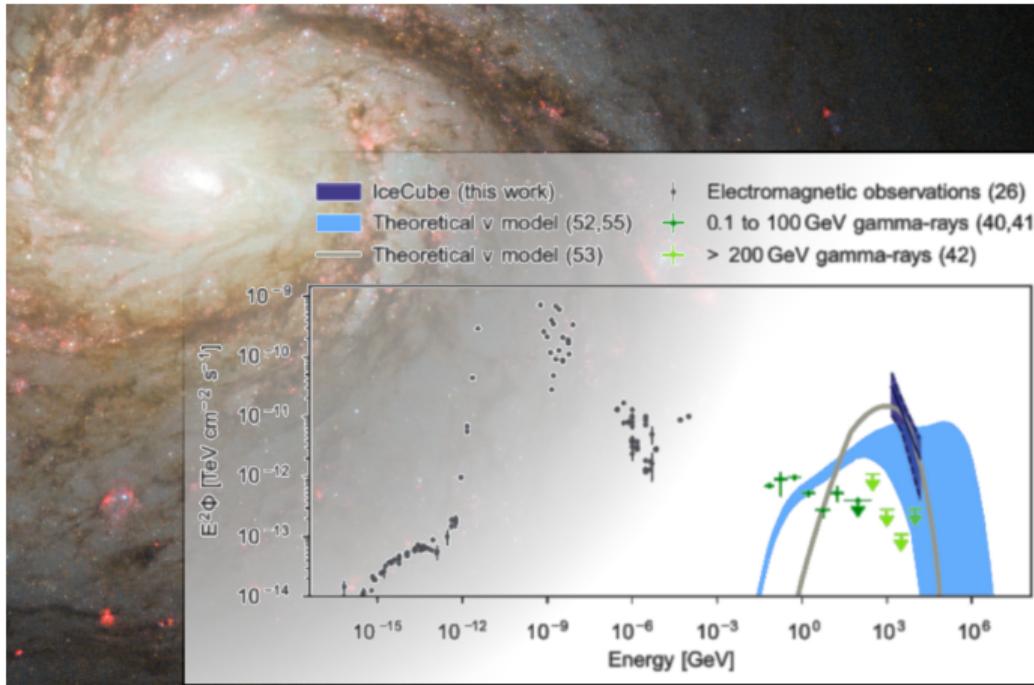


Figure: The Galaxy NGC 1068 has been identified as the source of the neutrino flux (*IceCube Collaboration, 2022, Science, 378, 6619*).

DANIELE VILLA 10/05/2023

Investigation methods

System spans several orders of magnitude:

- The Accretion Disk can extend for **thousands of $r_g = \frac{GM}{c^2}$** away from the Black Hole
- **Turbulence** develops on sub- r_g scales
- Recent efforts (*B. Ripperda et al., 2019, The Astrophysical Journal Supplement Series, 244, 1 – B. Ripperda et al., 2022, The Astrophysical Journal Letters, 924, 2.*) even attempt to **resolve the resistive scale** near the Black Hole → magnetic reconnection

The system is described using **General Relativistic MHD (GRMHD)**:

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The system is described using **General Relativistic MHD (GRMHD)**:

$$\nabla_\mu(\rho u^\mu) = 0 \quad (7)$$

$$\nabla_\mu(T^\mu_\nu) = 0 \quad (8)$$

$$\nabla_\mu({}^*F^\mu_\nu) = 0 \quad (9)$$

Accretion Disk Evolution

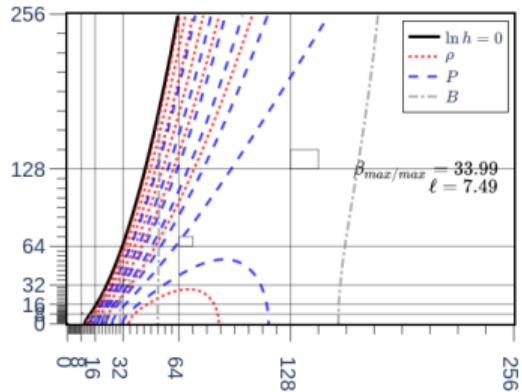


Figure: (L.G. Fishbone and V. Moncrief, 1976, *The Astrophysical Journal*, 207.)

Initial condition is **hydrodynamic** solution → Add “dynamically negligible” poloidal magnetic field → Accreted magnetic flux determines regime (MAD vs. SANE) (Y. Mizuno, 2022, *Universe*, 8, 2.).

Accretion Disk Evolution

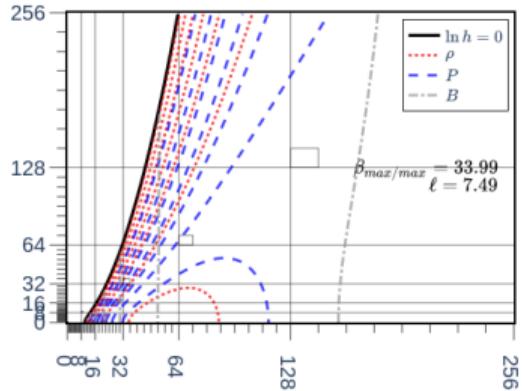


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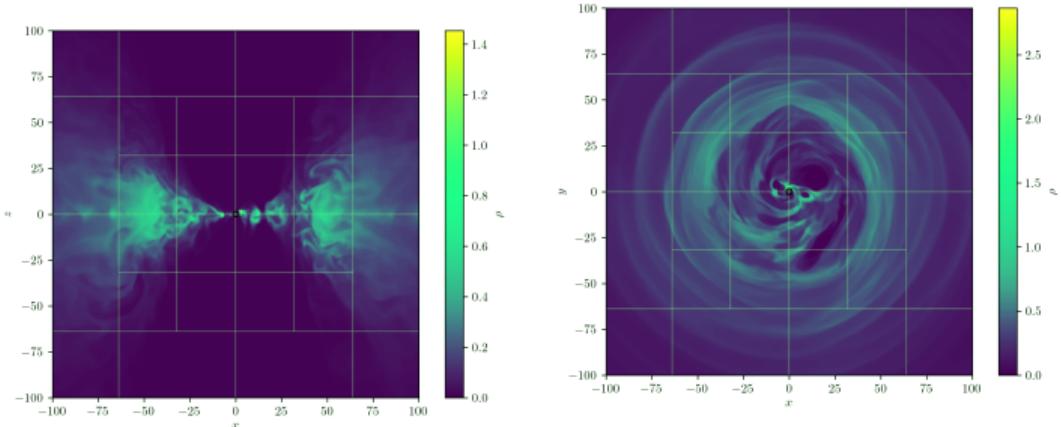


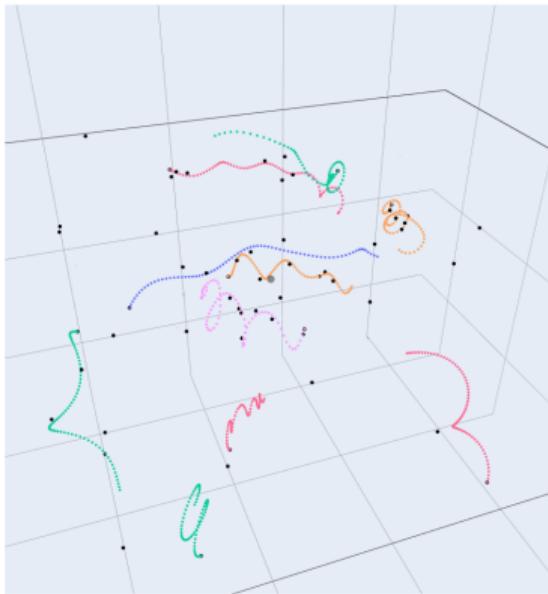
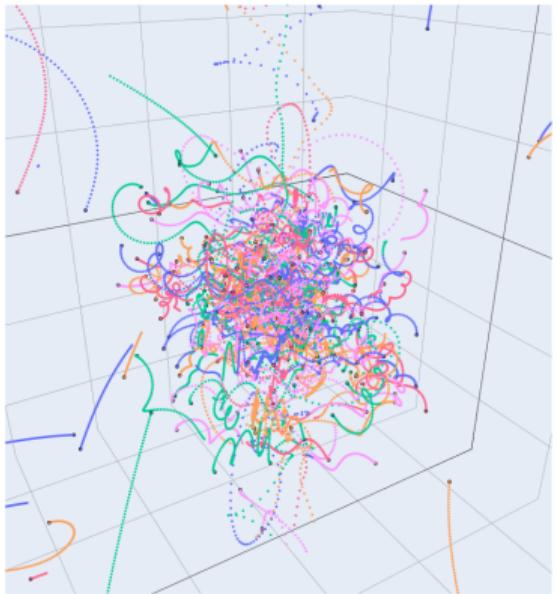
Figure: Density of the magnetized fluid around the Black Hole. View in the equatorial plane.

Figure: Density of the magnetized fluid around the Black Hole. View perpendicular to the equatorial plane.

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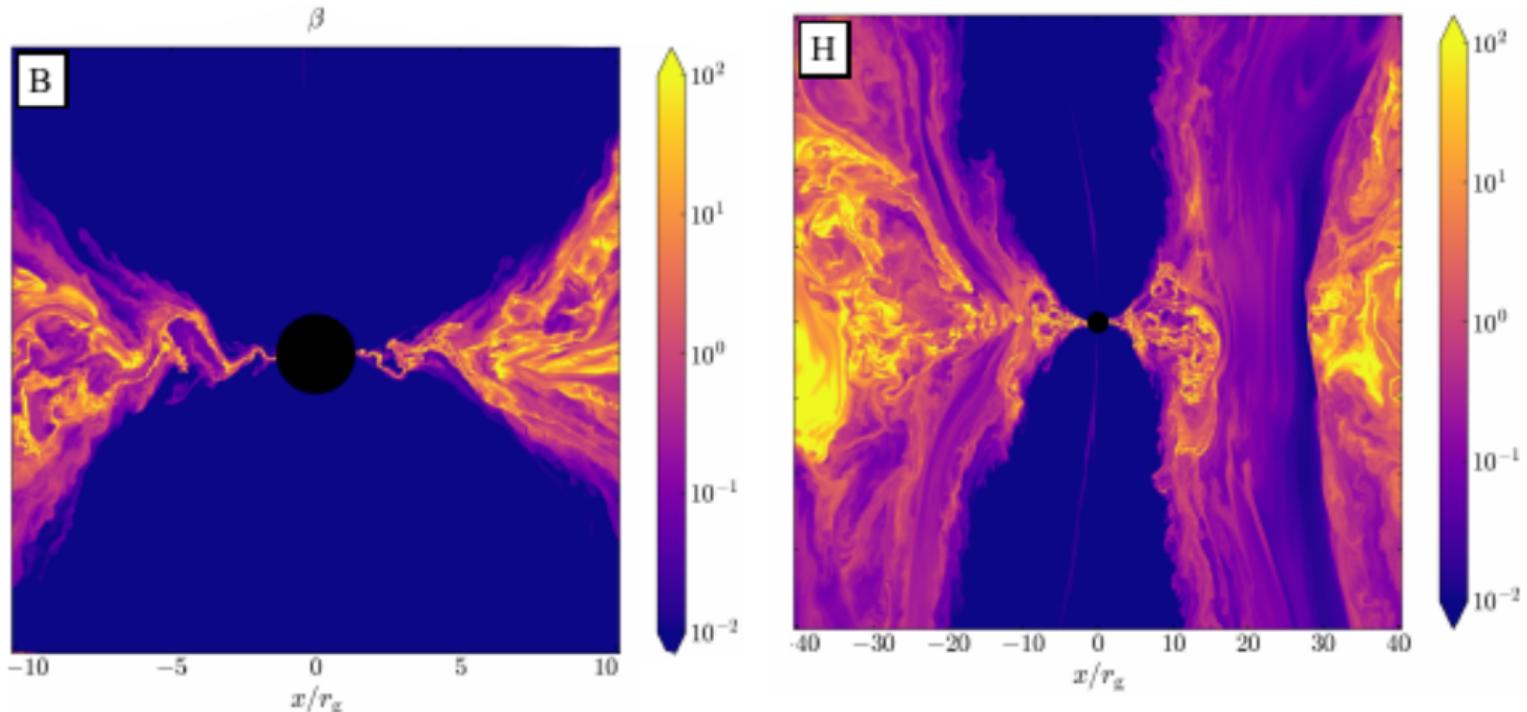


Adding particles



Particle trajectories in a MAD accretion disk.

Accretion Disk Evolution



Figures from Ripperda et al., 2022: resistive GRMHD

AthenaK Performance Scaling

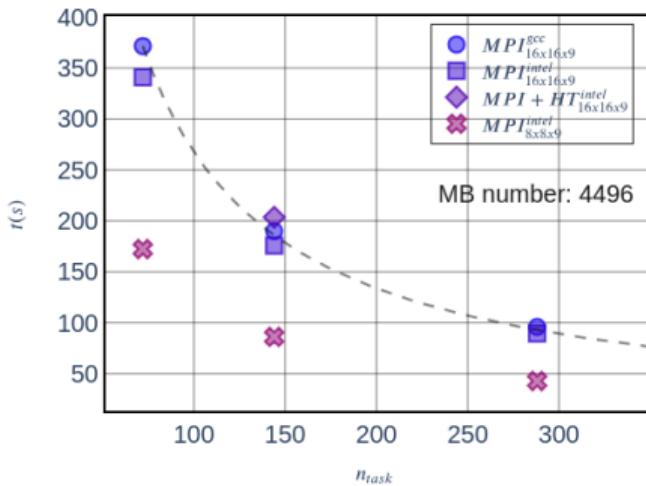


Figure: Performance scaling on CPU

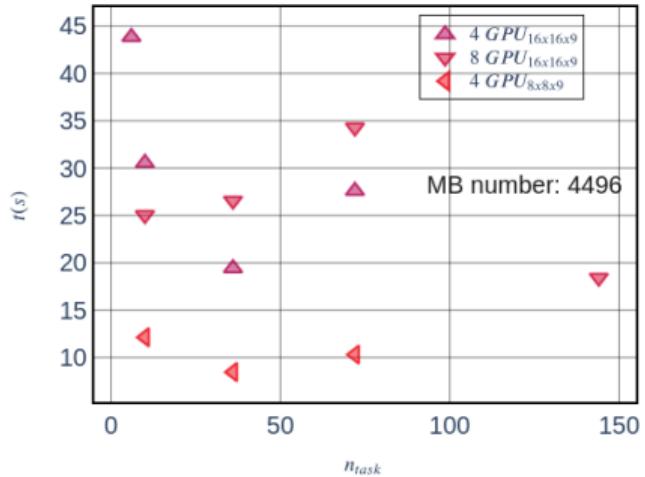


Figure: Performance scaling on GPU

AthenaK Performance Scaling

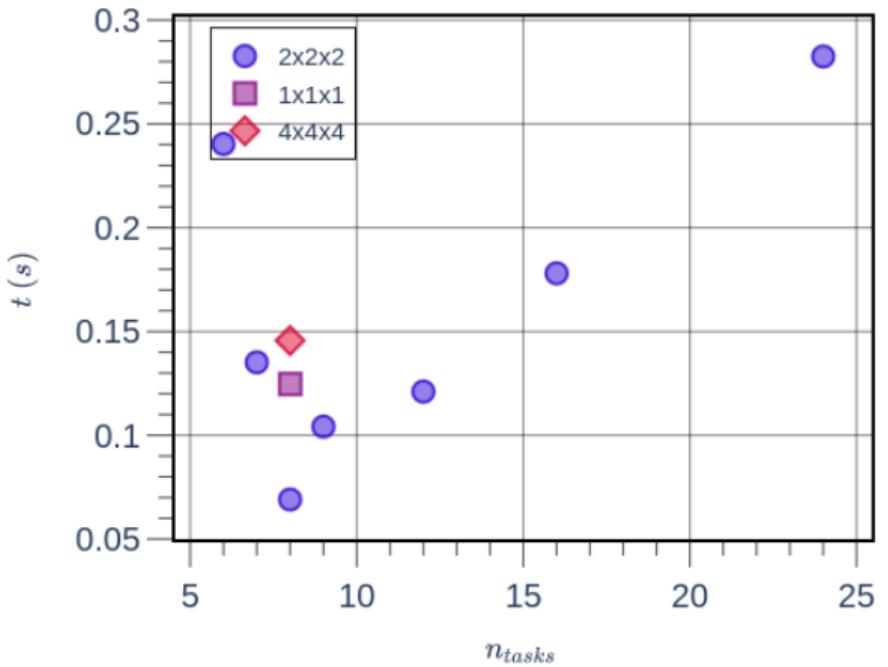


Figure: Performance scaling on GPU depending on domain setup.

Validation of particles in AthenaK

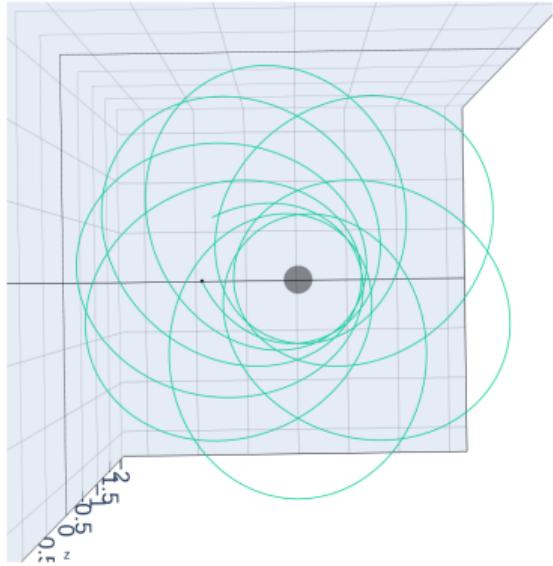


Figure: Precessing orbits in Schwarzschild (non-rotating black hole) spacetime.

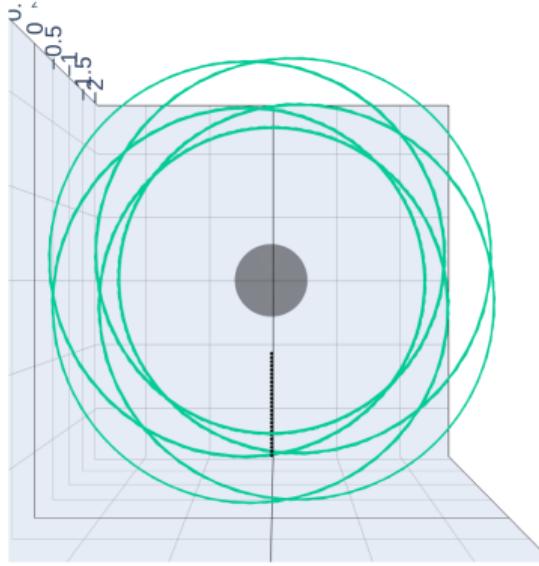


Figure: Closed orbits in Kerr (rotating black hole) spacetime.

Energy conservation properties of the algorithm

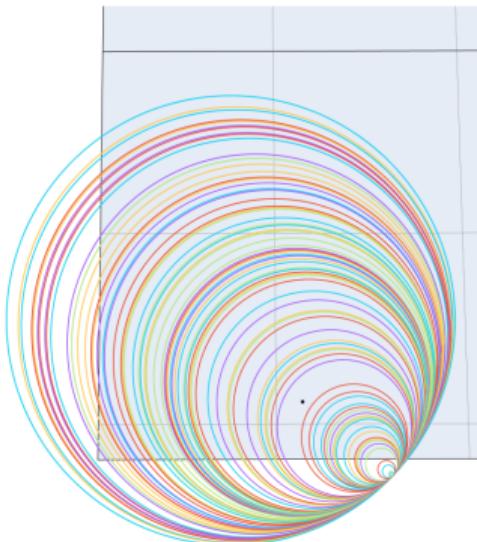


Figure: Larmor orbits in uniform magnetic field.

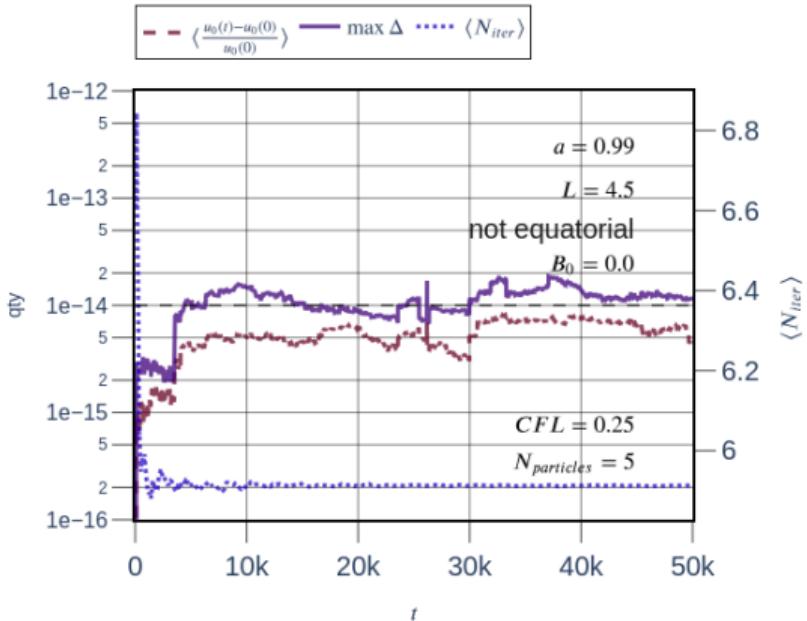


Figure: Energy conservation of the Hamiltonian scheme for gravitational dynamics.



Accretion Disk Evolution

