

#### **Experimental observation of Petschek's double shock structure in a two-fluid magnetic reconnection layer**

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Fast reconnection model conceived by H. Petschek based on MHD in 1964 has been very popular because it agrees well observations.

But its exact mechanism has not been well understood.



Petschek Model (1963) North Sato & Hayashi (1985)

**Uzdensky and Kulsrud (2000)**

## Outline

- Petshek Model (1964)
- Two-fluid physics of a typical magnetic reconnection layer
- The presence of sharp electron current layers at the separatrices
- **Comparison of our results with space data and simulations**
- Discussions and summary

### **=> A progress report for quantitative model of energy conversion and inventory**

• Petschek model has been studied in the MRX

Advanced simulation for the MRX reconnection layer in the two-fluid regime (J. Yoo)



Figure 3.4. Petschek model. [From Yamada et al. (2010).]



# Dynamics of 2-fluid reconnection layer



### Characteristics of MRX Data: Yamada et al, Nature Comms (2014), Princeton University Press (2022)





#### **Our recent simulation agrees very well with the MMS data and Parker Probe data**

### We have focused our study at the separatrix areas



#### **Our most recent MRX results verified a strong current sheet at the separatrix for Z > 2 di**



Most recent MRX data at  $z = 3 d_i$ 

## **2-D j<sub>e</sub>(R, Z) profiles support the Petshek-like feature**



Most recent MX data

#### **2-D V<sub>iz</sub> profiles support the Petshek-like feature**



## • **Petschek concept has been revived in MRX data**



More accurate measurement will be carried out in our FLARE device for a larger system size

$$
L>10\;di
$$

# Summary and future goals

- Petshek's double shock (1963) is observed on collisionless reconnection in MRX
- The presence of sharp electron current layers at the separatrices
- Our results are explained by two-fluid physics
- Ion acceleration and heating by re-magnetization
- **Our results agree well with space data**
- **More accurate measurements will be carried out in FLARE deice**
- **A simple model is being developed for energy conversion and partitioning**

**=> Quantitative model of energy conversion and inventory**

**A simple analytical model has been developed based on electron force balance:**

$$
E_R \approx V_{ey} B_Z - \frac{1}{en_e} \frac{\partial p_e}{\partial R}
$$
 (1)

We calculate the potential profile using Harris sheet expression (Yamada et al, PoP, 2000)

$$
B_z = -B_0 \tanh\left(\frac{x}{\delta}\right)
$$

$$
j_y = \frac{B_0}{\mu_0 \delta} \operatorname{sech}^2\left(\frac{x}{\delta}\right)
$$

sheet thickness  $\delta$  is given by

$$
\delta = \frac{c}{\omega_{\text{pi}}}\frac{\sqrt{2(T_e + T_i)/m_i}}{V_i - V_e} = \frac{c}{\omega_{\text{pi}}}\frac{\sqrt{2}V_s}{V_{\text{drift}}}
$$

We obtain an analytical expression of potential  $(\Phi)$  profile

$$
\Phi_p(x) = -\Phi_0 \operatorname{sech}(x) \quad \text{with} \quad \Phi_0 = \frac{B_0^2}{2\mu_0 e} - \frac{\Delta T_e}{e}.
$$

**A half of incoming Poynting flux energy**   $\Rightarrow$  Potential Energy:  $e\Phi_0 = (1/2) (B_0^2/\mu_0)$ 



### 2D PIC simulation on energetics; by J. Jara Almonte and W. Daughton



The energy partitioning does not strongly depend on the size of monitoring boundary



- Ion flux : e-flux  $\sim$  2:1
- Enthalpy flux dominates
- Magnetic energy outflow substantial

It is notable that energy deposition to ions is generally larger than to electrons. **Since the electrons' heat transport loss is larger than ions',** 

 $\equiv$   $\sum_i$   $\gg$   $\sum_e$ 

 $\Leftrightarrow$  Combination of larger energy deposition and better confinement for ions