



Electrostatic waves and electron holes in PIC simulations of the Earth's bow shock.



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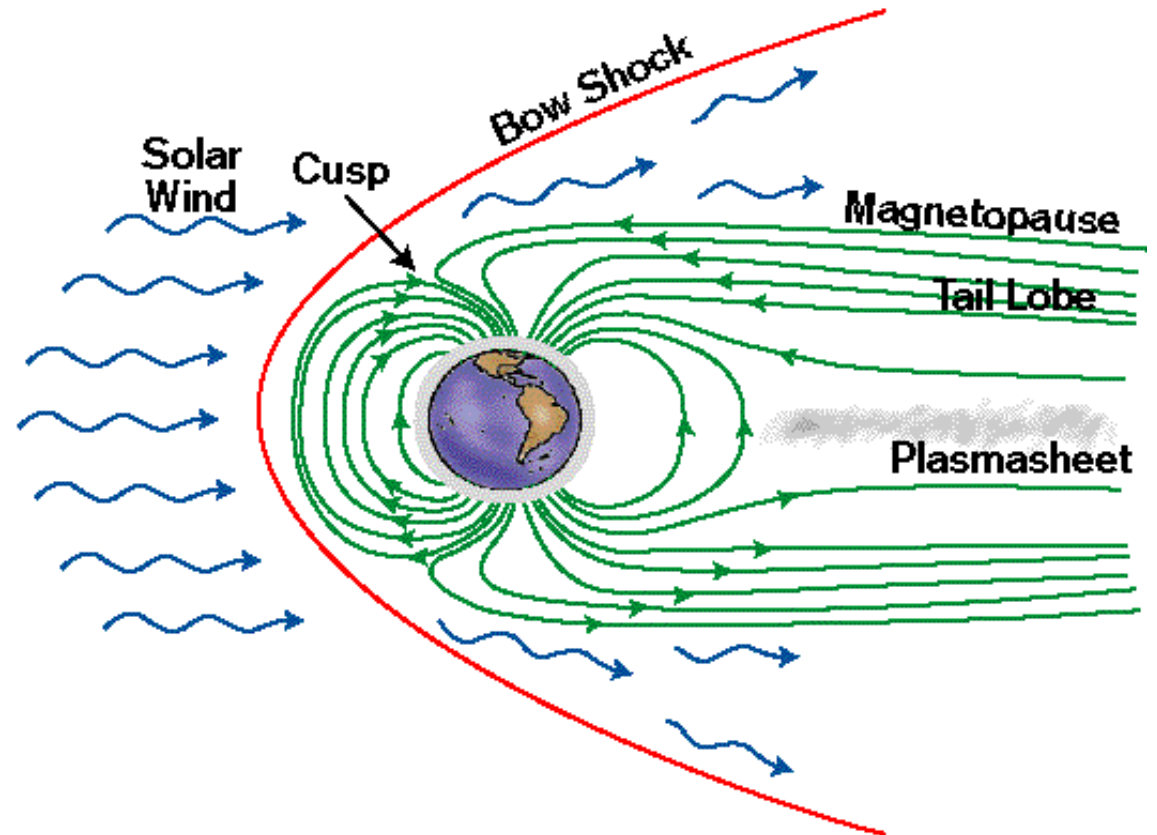
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The Earth's Bow Shock

- Nonrelativistic shocks: $v_{sh} \ll c$
- Sonic Mach number: $M_S = v_{sh}/c_s \approx 2-10$
- Alfvén Mach number: $M_A = v_{sh}/v_A \approx 2-10$

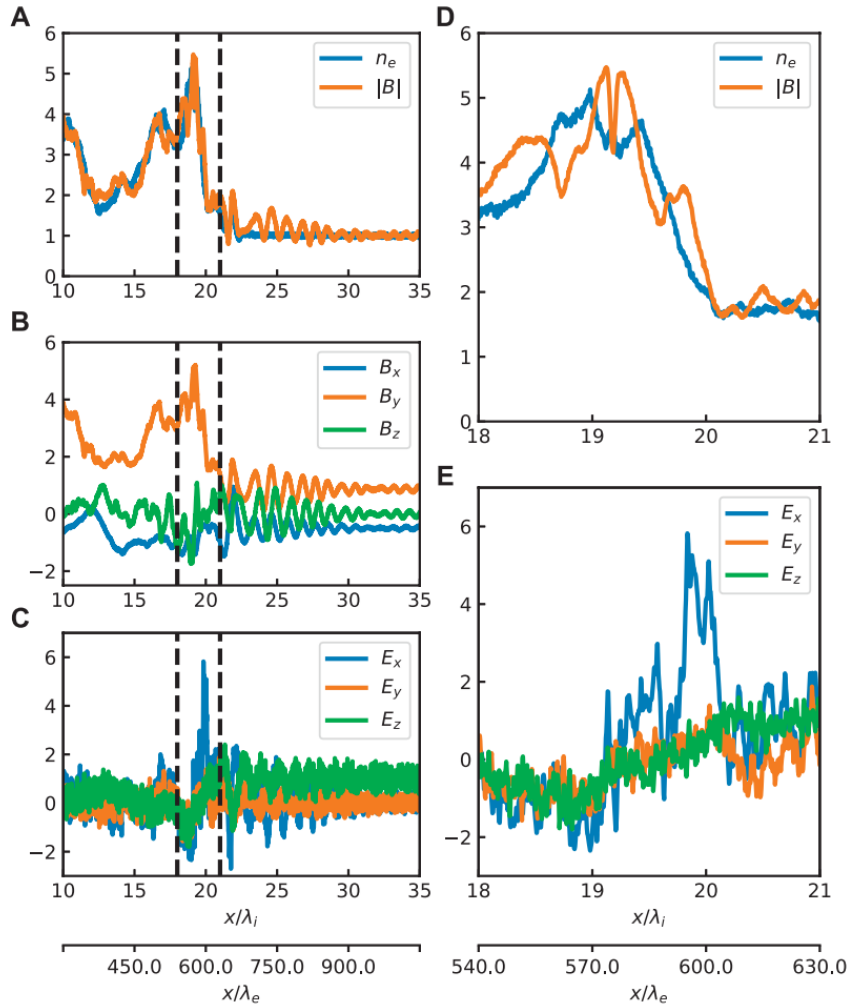


PIC simulations vs in-situ measurement: electrostatic waves

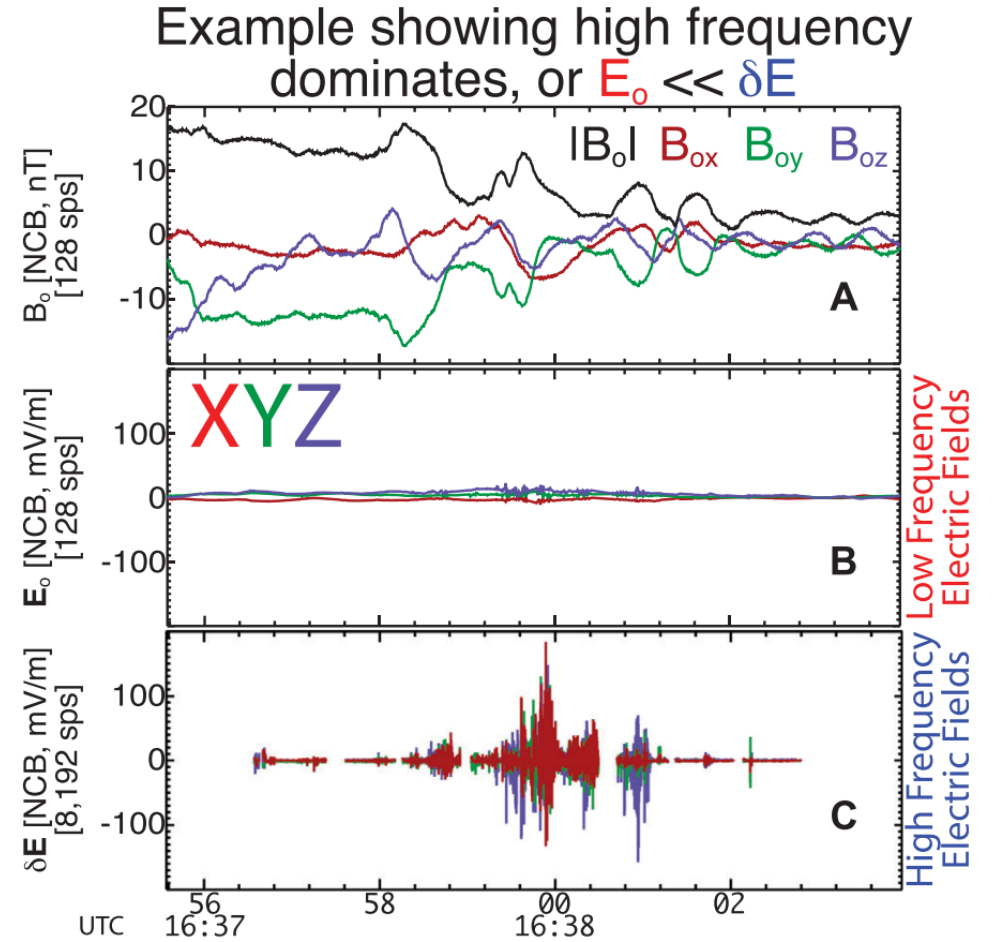
(Wilson et al. 2021)



PIC simulations



In-situ



Electrostatic waves at/near collisionless shocks (in-situ)



(Wilson et al. 2021)

Wave Name	Polarization or waveform	Frequency ^α and/or Appearance	Scale Length ^β	Free energy source or wave source
LHW	linear ⊥ to \mathbf{B}_o or oblique to \mathbf{B}_o	$f_{sc} \sim 5\text{--}40$ Hz $f_{sc} \lesssim f_{lh}$ symmetric modulated sine waves ^p	$k \lambda_e \lesssim 1$	currents ^k , density gradients ^λ , Electron heat flux ^σ , or MTSI ^θ
IAW	linear to \mathbf{B}_o	$f_{sc} \sim 10^2\text{--}10^4$ Hz $f_{rest} \lesssim f_{pi}$ symmetric ⁿ modulated sine waves	$\lambda \geq 2\pi\lambda_{De}$	currents ^δ , gyrating/reflected ions ^ζ , or electron heat flux ^ξ
ECDI	elliptical or “Tear-drop”- shaped oblique to \mathbf{B}_o	$f_{sc} \sim 10^2\text{--}10^4$ Hz $f_{rest} \sim \text{mix}^e$ asymmetric ⁿ modulated sine waves	$k \lambda_e \lesssim 1$ and $k \lambda_{De} \lesssim 1$	relative drift between incident electrons and reflected ions ^δ
ESW	bipolar pulse to \mathbf{B}_o else unipolar	$f_{sc}^{-1} \sim \text{few } 10 \text{ s of ms}$ isolated or trains of pulses	$\lambda \geq \lambda_{De}$	electron beams ^δ or nonlinear wave decay ^δ
LW ^ψ	linear to \mathbf{B}_o or elliptical	$f_{sc} \sim 10\text{--}60$ kHz symmetric modulated	$k \lambda_e \lesssim 1^u$	electron beams ^x and/or nonlinear wave decay ^v

Properties of IAW and ESW: $\delta E / E_0 > 50$

$$\lambda_{ESW} \approx 10\lambda_D \approx 0.05\lambda_{se}$$

Electrostatic waves at/near collisionless shocks (PIC simulations)



- **Buneman waves** at the shock foot of quasi-perpendicular high (Shimada & Hoshino, 2000; Hoshino & Shimada, 2002; Amano & Hoshino, 2007, 2009; Bohdan et al., 2017, 2019a, 2019b) and low (Umeda et al., 2009) Mach number shocks.
- **Electron Bernstein mode** (Muschiatti & Lembège, 2006; Yu et al., 2022) can be excited in moderate Mach number perpendicular shocks.
- **Ion-acoustic waves** can be driven by the drift motion of preheated incoming ions relative to the decelerated electrons at the shock foot of high Mach number perpendicular shocks (Kato & Takabe, 2010b, 2010a).
- **Electron-acoustic waves** can be observed both in the shock foot as a result of the MTSI (Matsukiyo & Scholer, 2006) or in the electron foreshock of oblique shocks (Bohdan et al., 2022; Morris et al., 2022).
- **Electrostatic Langmuir waves** can be generated via the electron bump-on-tail instability at the foreshock region of oblique high-beta shocks (Kobzar et al., 2021)

Properties of electrostatic waves:

$$\delta E / E_0 \approx 1$$

$$\lambda_{EW} \approx \lambda_{se}$$

PIC simulations vs in-situ measurement



PIC simulations

$$\delta E / E_0 \approx 1$$

$$\lambda_{EW} \approx \lambda_{se}$$

In-situ

$$\delta E / E_0 > 50$$

$$\lambda_{ESW} \approx 10\lambda_D \approx 0.05\lambda_{se}$$

PIC simulations vs in-situ measurement



PIC simulations

$$\delta E / E_0 \approx 1$$

$$\lambda_{EW} \approx \lambda_{se}$$

In-situ

$$\delta E / E_0 > 50$$

$$\lambda_{ESW} \approx 10\lambda_D \approx 0.05\lambda_{se}$$

Simple explanation

Very often $\lambda \propto v_{sh}$ (two-stream instability)

Also $\delta E \propto v_{sh}$ (available energy). $E_0 = B_0 v_{sh} \propto v_{sh}^2$ (assuming constant M_A), therefore $\delta E / E_0 \propto v_{sh}^{-1}$

$$v_{sim} = 0.2c$$

$$\lambda_{EW} \approx \lambda_{se}$$

$$\delta E / E_0 \approx 1$$

$$v_{real} = 0.002c \approx 600 \text{ km/s}$$

$$\lambda_{ESW} \approx 0.01\lambda_{se}$$

$$\delta E / E_0 \approx 100$$

PIC simulations vs in-situ measurement



PIC simulations

$$\delta E/E_0 \approx 1$$

$$\lambda_{EW} \approx \lambda_{se}$$

In-situ

$$\delta E/E_0 > 50$$

$$\lambda_{ESW} \approx 10\lambda_D \approx 0.05\lambda_{se}$$

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$$v_{sim} = 0.2c$$

$$\lambda_{EW} \approx \lambda_{se}$$

$$\delta E/E_0 \approx 1$$

Does it work?

$$v_{real} = 0.002c \approx 600 \text{ km/s}$$

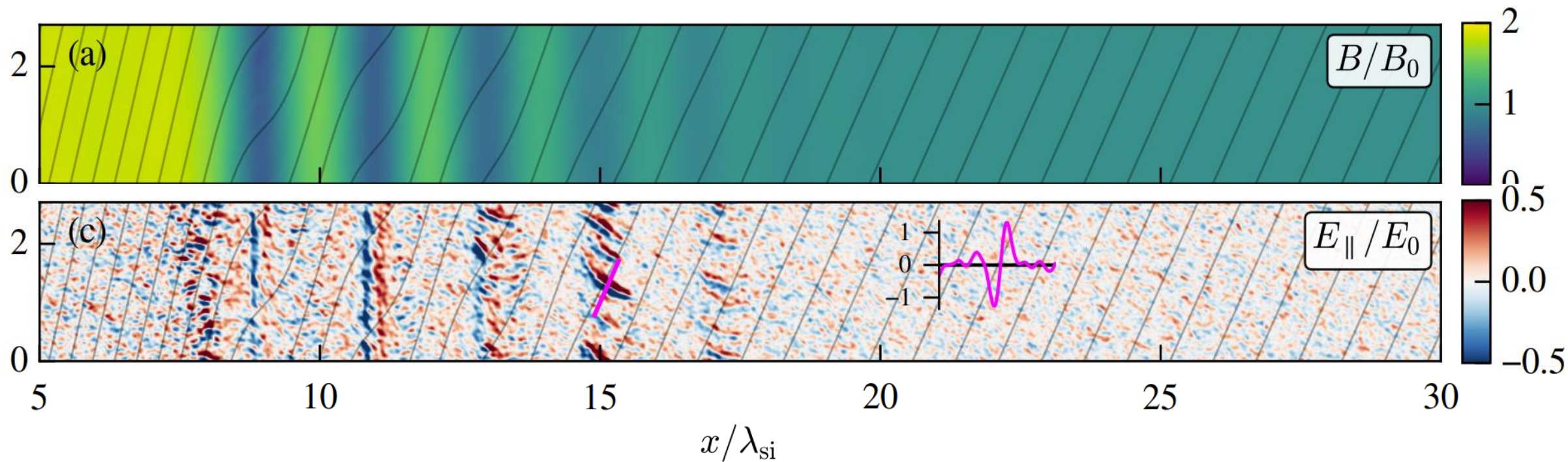
$$\lambda_{ESW} \approx 0.01\lambda_{se}$$

$$\delta E/E_0 \approx 100$$

Shock simulations

$$M_A = 1.8, M_S = 4, \Theta_{Bn} = 65^\circ$$

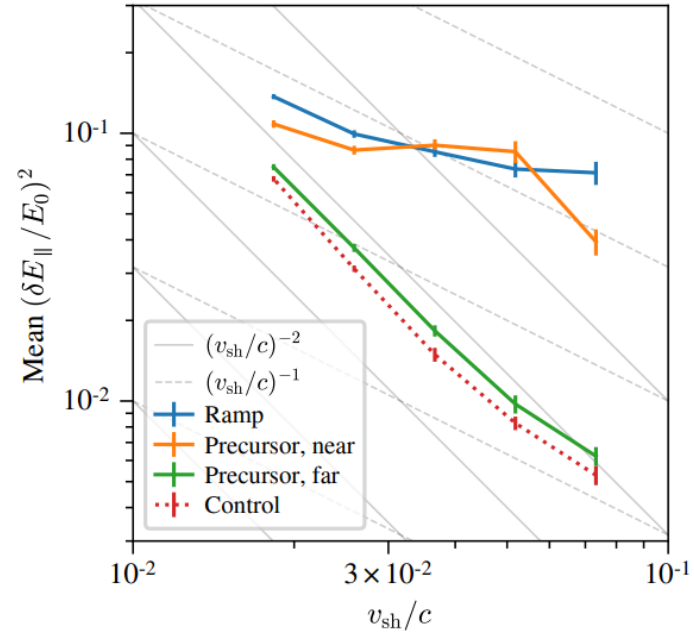
Run	m_i/m_e	v_{sh}/c^\dagger	v_0/c	Width (d_i)	Δx (d_e)	ω_{pe}/Ω_e
A	200	0.0733	0.0338	2.90	0.143	1.76
B	200	0.0518	0.0238	2.71	0.100	2.49
C	200	0.0366	0.0168	2.90	0.071	3.52
D	200	0.0259	0.0119	2.71	0.050	4.99
E	200	0.0183	0.0084	2.90	0.036	7.05



Wave parameters scaling

Electrostatic wave power

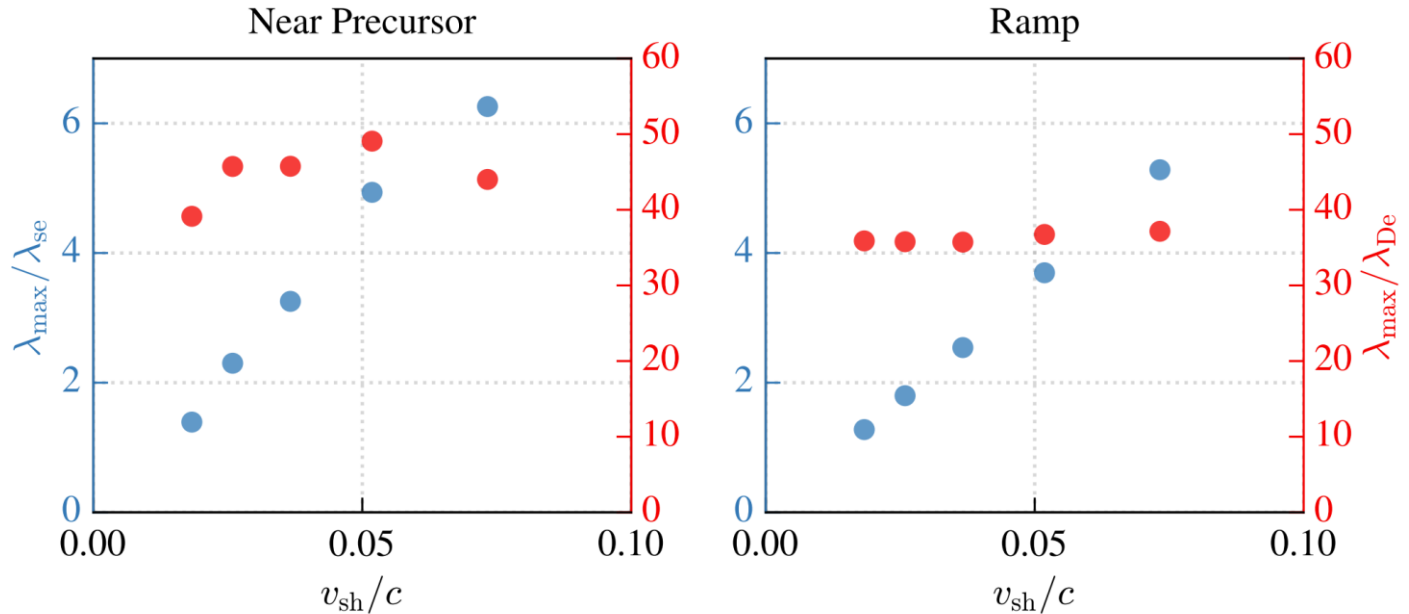
$$\delta E / E_0 \propto v_{sh}^{-0.5}$$



Wavelength

$$\lambda_{EW} / \lambda_{se} \propto v_{sh}$$

$$\lambda_{EW} / \lambda_D \propto v_{sh}^0$$

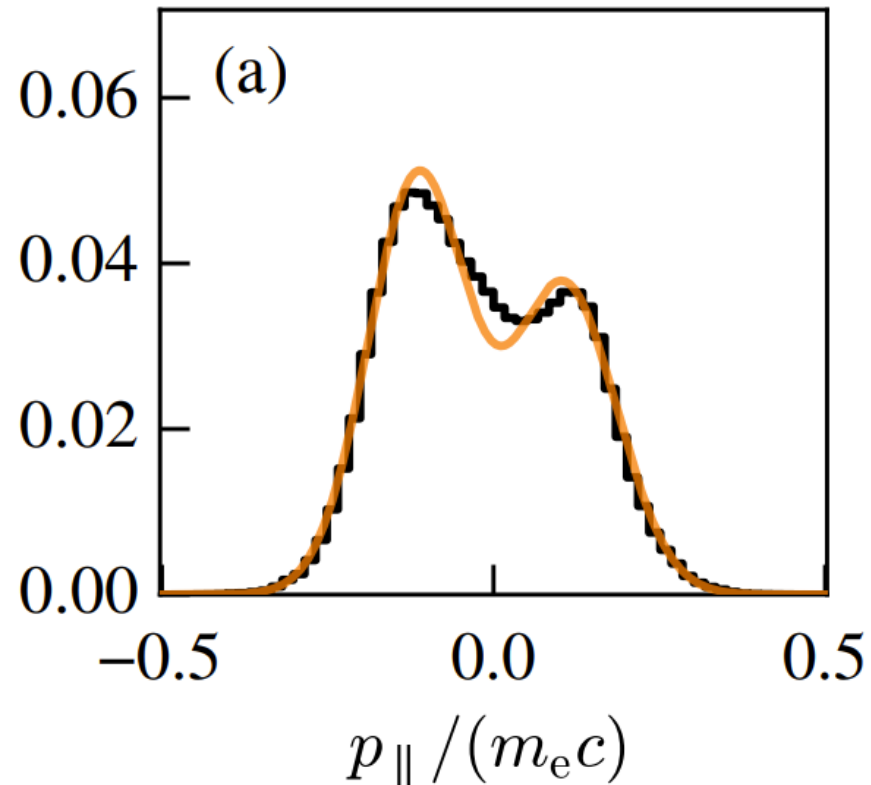


Linear dispersion analysis

Two-stream electrostatic instability



Electron distribution at the shock ramp (run B)



Parameters of the electron distribution

Run	$\frac{n_1}{n_2}$	$\frac{v_{dr}}{v_{sh}}$	$\frac{v_{dr}}{v_{th,1}}$	$\frac{v_{dr}}{v_{th,2}}$	v_{sh}/c
A	1.52	4.14	2.79	2.78	0.0733
B	1.36	4.15	2.88	2.75	0.0518
C	1.33	4.16	2.96	2.61	0.0366
D	1.28	4.19	2.99	2.45	0.0259
E	1.26	4.28	3.02	2.40	0.0183
S*	1.35	4.19	2.93	2.59	0.0010

*S – synthetic run, $v_{sh} \approx 312 \text{ km/s}$

Linear dispersion analysis (LDA)

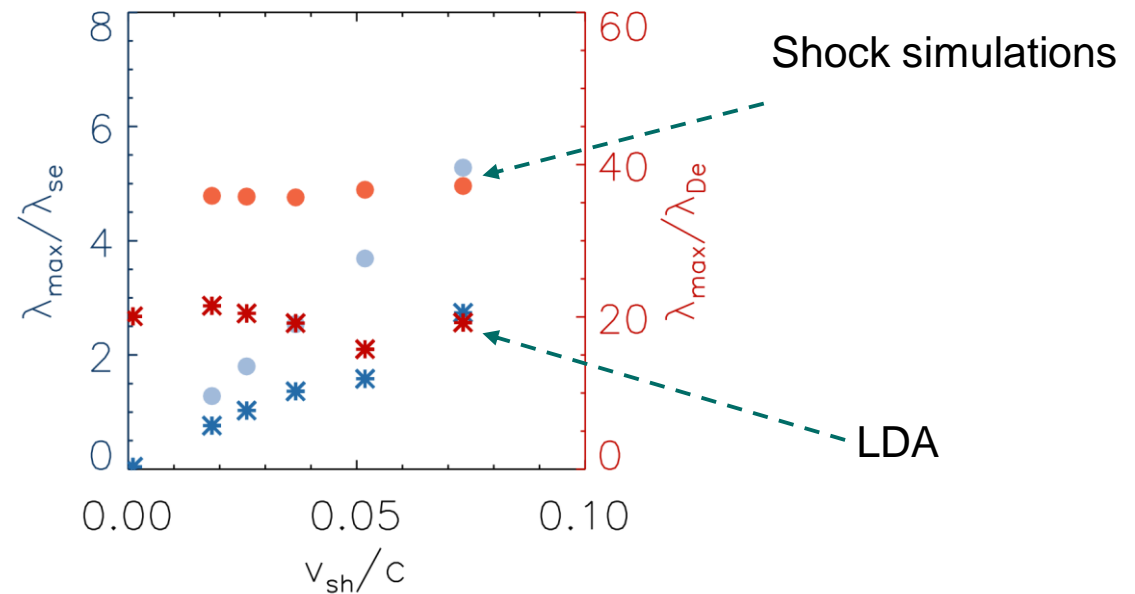
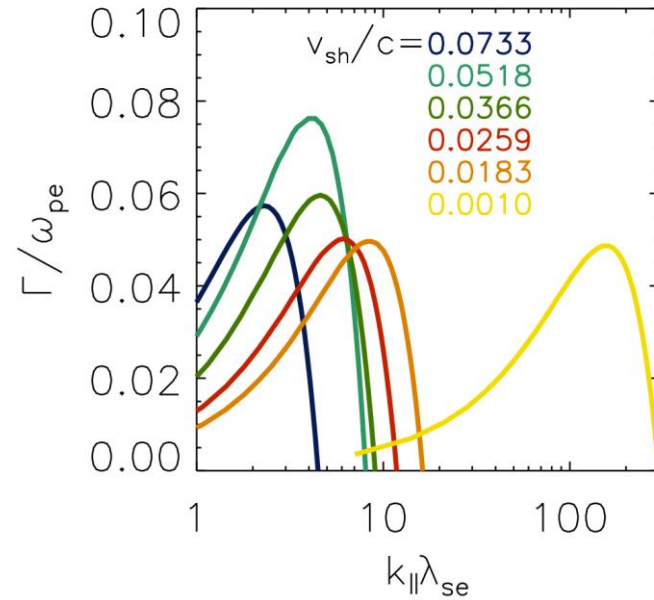
Electron-acoustic waves

Electrostatic waves are identified as electron-acoustic waves driven by two counterstreaming hot electron beams.

$$\lambda_{LDA}/\lambda_{se} \propto v_{sh}$$

$$\lambda_{LDA}/\lambda_D \propto v_{sh}^0$$

But $\lambda_{ES,sh}/\lambda_{LDA} \approx 2$



Periodic-boundary-condition simulations

Setup

\longrightarrow $v_1 = \frac{v_{dr}}{2}$ $v_{th,1} = \frac{v_{dr}}{3}$	\longleftarrow $v_2 = -\frac{v_{dr}}{2}$ $v_{th,2} = \frac{v_{dr}}{5}$
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Scanned parameters

$$v_{dr} \approx (0.08 - 0.3)c$$

$$\frac{m_i}{m_e} = 200 - 1836$$

$$N_{ppc} = 40 - 2560$$

Ion content

Reference PBCS

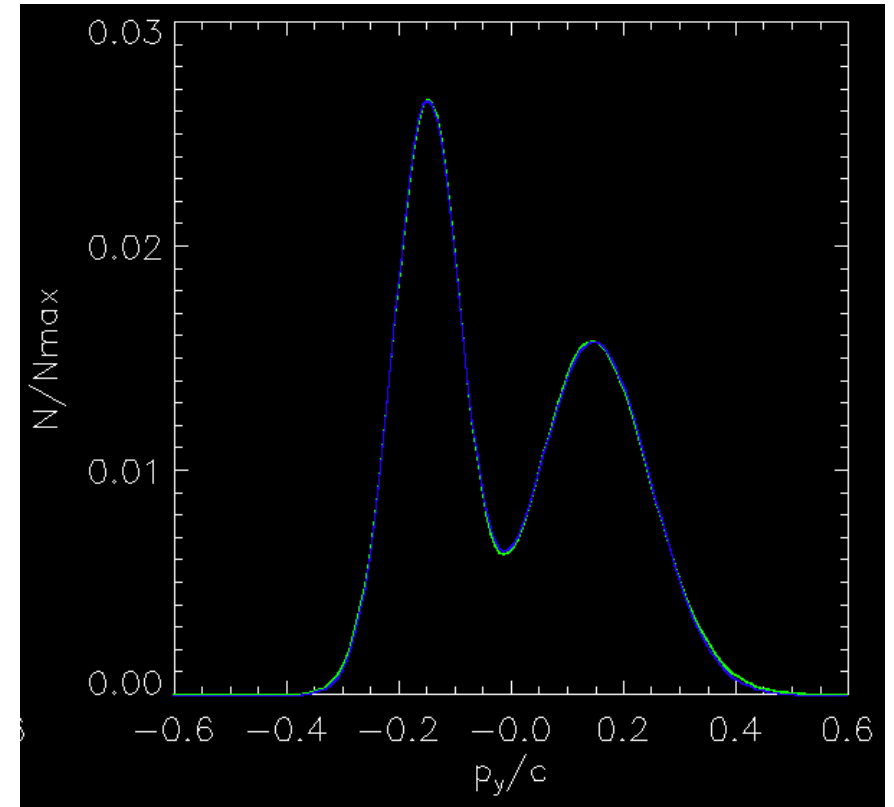
$$v_{dr} = 4v_{sh, run A}$$

$$n_1 = n_2$$

$$N_{ppc} = 2560$$

$$v_{th,1} = \frac{v_{dr}}{3}$$

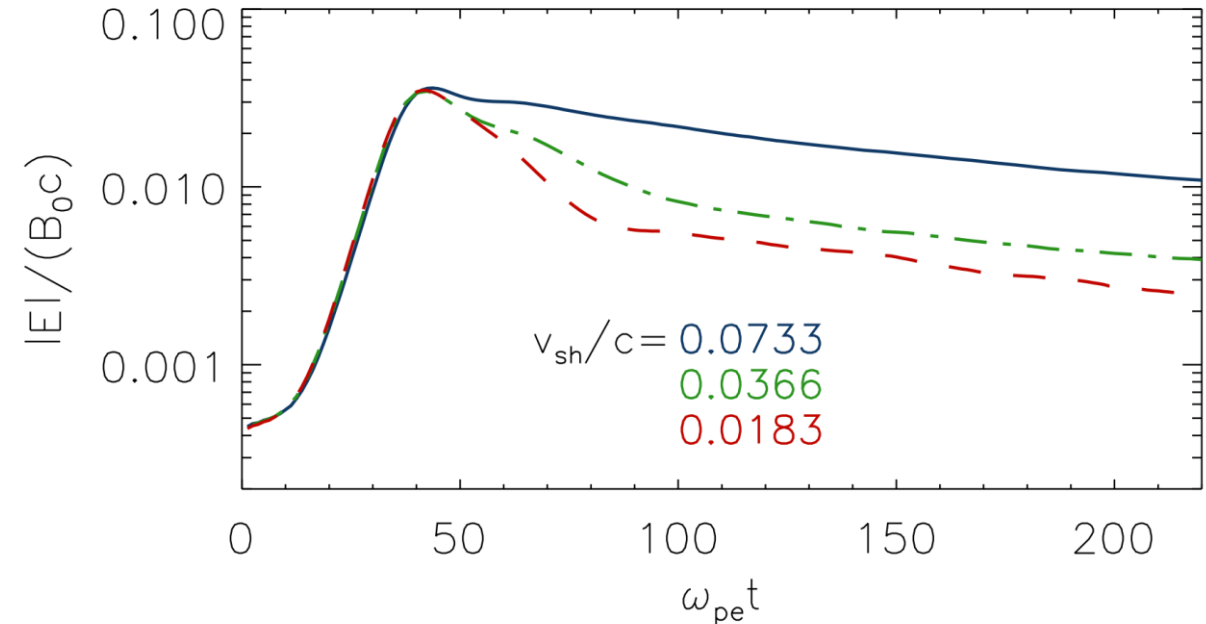
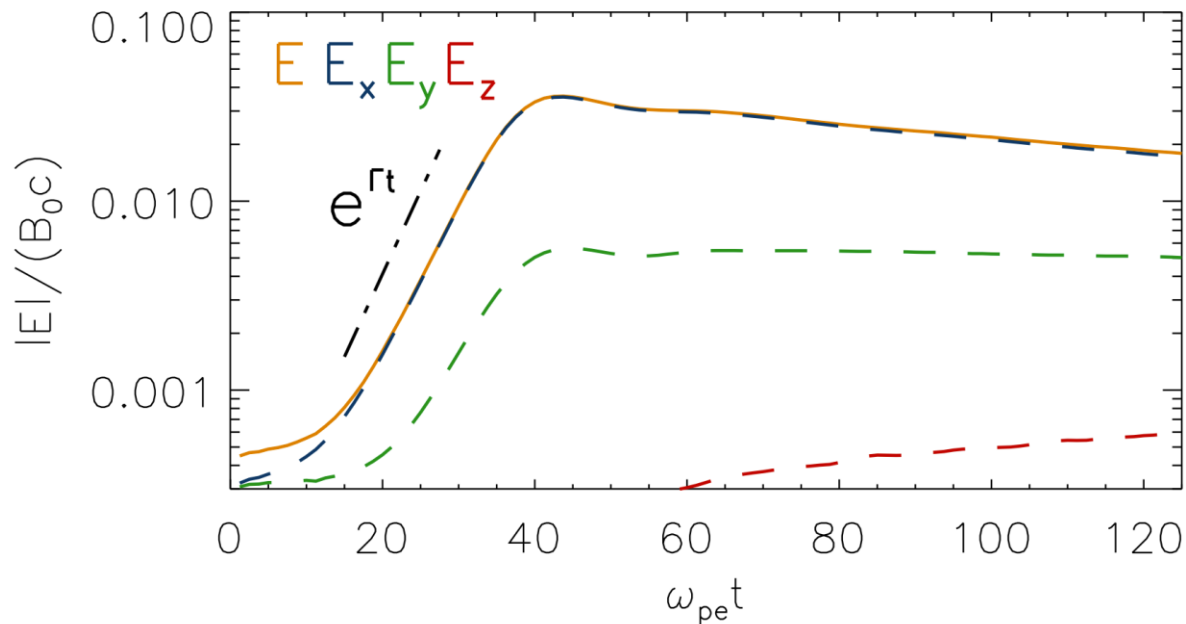
$$v_{th,2} = \frac{v_{dr}}{5}$$



Periodic-boundary-condition simulations

Results

- Good match between PBCS and linear dispersion analysis ($\Gamma/\omega_{pe} \approx 0.185$, the prediction is 0.2)
- $\text{Max}(\delta E/(B_0 c))$ does not depend on v_{sh} , therefore $\delta E/E_0 \propto v_{sh}^{-1}$
- At $\omega_{pe} t > 50$, $\delta E/(B_0 c) \propto v_{sh}^{-1}$, therefore $\delta E/E_0 \propto v_{sh}^0$

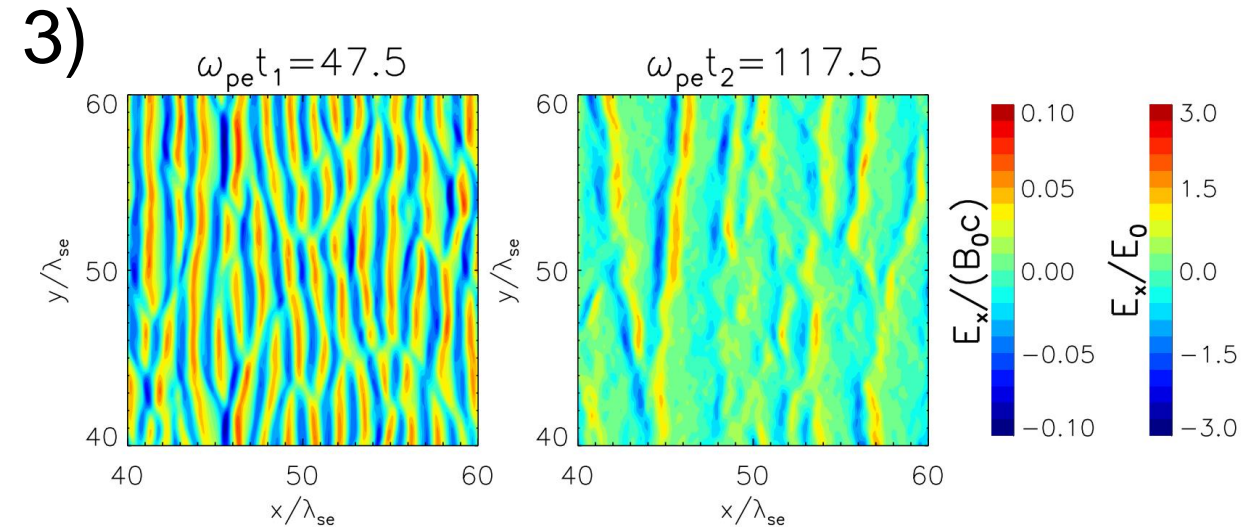
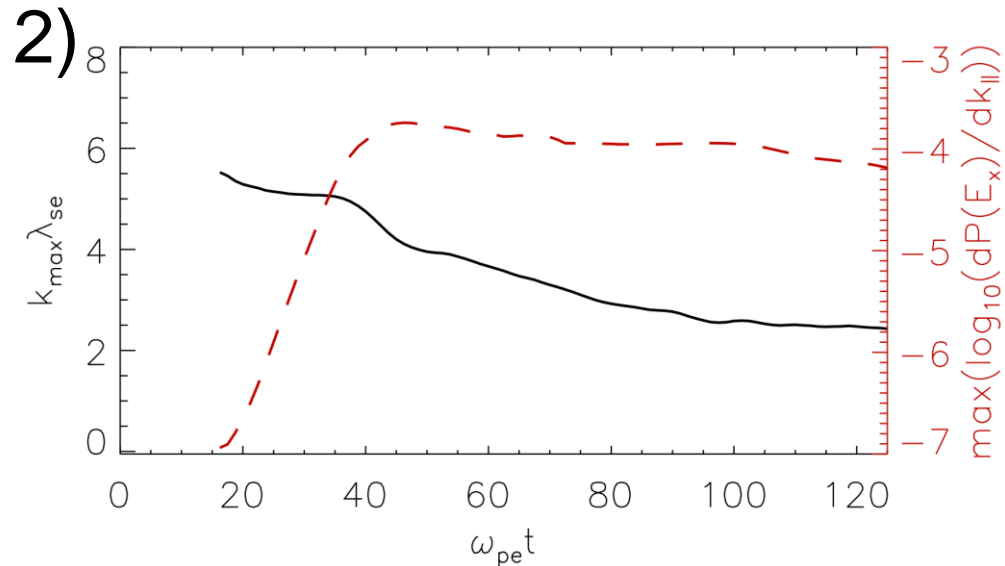
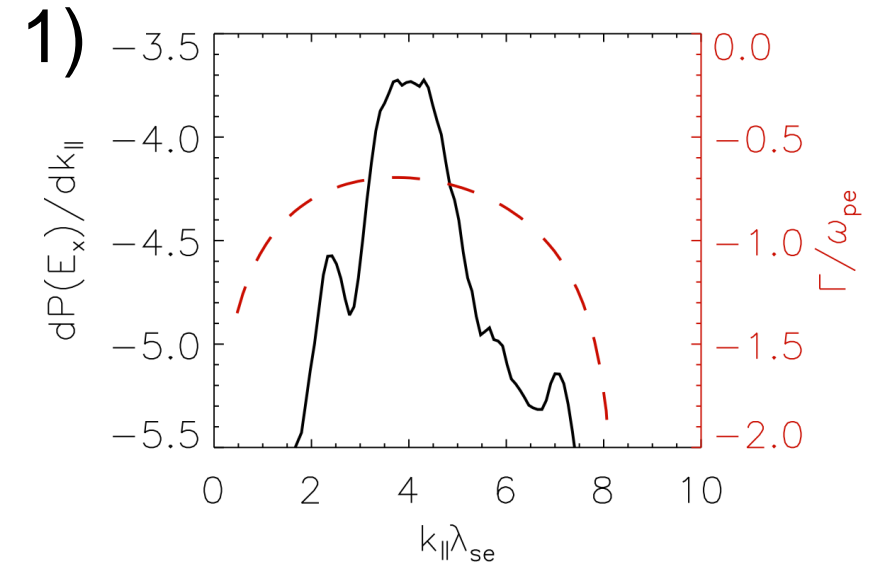


Periodic-boundary-condition simulations



Results

- 1) The wavelength is consistent with predictions and $\lambda_{ES}/\lambda_{se} \propto v_{sh}$
- 2) The wavelength $\lambda_{ES}/\lambda_{se}$ is decreasing with time
- 3) Structure becomes closer to solitary waves at later stages of EAW development



Conclusions



- Driving conditions for EAWs are independent on the shock velocity.
- The amplitude scales as $\delta E/E_0 \propto v_{sh}^{-0.5}$ or $\delta E/E_0 \propto v_{sh}^{-1}$, therefore $\delta E/E_0$ in real shocks could be larger than **100**.
- The wavelength scales as $\lambda_{ESW}/\lambda_D \propto v_{sh}^0$, therefore $\lambda_{ESW} \approx 40 \lambda_D \approx 300m$.

arXiv > physics > arXiv:2408.01699

Physics > Space Physics

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Electrostatic Waves and Electron Holes in Simulations of Low-Mach Quasi-Perpendicular Shocks

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Thank you

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