# **The problem with the r-modes of neutron stars**

N. Andersson and F. Gittins (2023). Astrophys. J. 945.2, 139. DOI: 10.3847/1538-4357/ [acbc1e](https://doi.org/10.3847/1538-4357/acbc1e)

F. Gittins and N. Andersson (2023). Mon. Not. R. Astron. Soc. 521.2, 3043. DOI: [10.1093/](https://doi.org/10.1093/mnras/stad672) [mnras/stad672](https://doi.org/10.1093/mnras/stad672)

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"Continuous gravitational waves and neutron stars workshop", Hannover, Germany, 19th Jun. 2024

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## **The r-modes**

- The r-modes are a family of oscillations that arise due to stellar rotation.
- They couple weakly to density perturbations and have (leading-order) frequency

$$
\omega_0 = \frac{2m}{l(l+1)}\Omega.
$$

An  $(l > 2)$  *r*-mode travels *retrograde* to the rotation in the co-rotating frame, but prograde in the inertial frame.



**Figure 1:** Streamlines of the  $l = 2$  r-modes with  $m = 1$  (left panel) and  $m = 2$  (right panel) in the co-rotating frame.

- The oscillations are generically unstable to gravitational radiation (Andersson [1998;](#page-12-0) Friedman and Morsink [1998\)](#page-12-1).
- Real neutron stars possess bulk and shear viscosities that compete with the growth of the instability. For some combination of the star's temperature and spin, the instability wins out.





- The instability may limit the rotation rates of young pulsars (Lindblom, Owen and Morsink [1998\)](#page-12-3) and mature accreting systems (Bildsten [1998\)](#page-12-4).
- Furthermore, a recent search has been conducted by the LIGO-Virgo-KAGRA Collaboration on the glitching PSR J0537-6910 (Abbott et al. [2021\)](#page-12-5).
- However, our present understanding of the r-modes relies on (i) Newtonian gravity and (ii) simple matter models.
- The situation is unsolved in relativity, since we (surprisingly) find a singular-eigenvalue problem (Kojima [1998\)](#page-12-6):

$$
(\varpi - \mu) \left[ \frac{1}{jr^4} \frac{d}{dr} \left( j r^4 \frac{d\Phi}{dr} \right) - v \Phi \right] = q \Phi.
$$



- This suggests that there is an issue with our assumptions.
- These problems need to be addressed to inform gravitational-wave searches and reliably extract neutron-star parameters from the signal.
- In this work, we extend the r-mode calculations to realistic nuclear-matter models.

• Mature neutron stars are cold and stratified by composition gradients,

$$
d\varepsilon = \frac{\varepsilon + p}{n_{\rm b}}dn_{\rm b} + n_{\rm b}\mu_{\Delta}dY_{\rm e} \quad \Longrightarrow \quad \varepsilon = \varepsilon(n_{\rm b}, Y_{\rm e}).
$$

• When the star is in  $\beta$  equilibrium,  $\mu_{\Delta} = \mu_{p} + \mu_{e} - \mu_{n} = 0$ , the matter is barotropic,  $\varepsilon = \varepsilon(n_{\rm b})$ , and possesses

$$
\Gamma = \frac{d \ln p}{d \ln n_{\rm b}}.
$$

• Deviations from equilibrium are encoded in the adiabatic index

$$
\Gamma_1 = \left(\frac{\partial \ln p}{\partial \ln n_{\mathsf{b}}}\right)_{Y_{\mathsf{e}}}
$$

*.*

#### **Simple matter models**

Common assumptions in the literature include (i)  $\Gamma_1 = \Gamma$  (neglecting nuclear reactions) and (ii)  $\Gamma_1 = \text{const} \neq \Gamma$ .

### **Realistic nuclear matter: stratification**



**Figure 3:** Stratification in nuclear-matter equations of state BSk19 and BSk21.

• A real neutron star is expected to be predominantly non-barotropic, with locally barotropic layers. Thus,  $\Gamma_1 \neq \text{const.}$ 

#### **Opportunity**

Probing the dense nuclear-matter equation of state is one of the main science goals of gravitational-wave astronomy.

- Consider a barotropic  $(\Gamma_1 = \Gamma)$  star.
	- When it is non-rotating, there exists time-independent solutions that do not source density perturbations.
	- These are zero-frequency (polar) g-modes—that have no stratification support—and (axial) r-modes.
	- As the star spins up, these solutions oscillate and mix, forming inertial modes.
	- For each  $l = |m|$ , one purely axial *r*-mode persists.
- Next, focus on a non-barotropic  $(\Gamma_1 \neq \Gamma)$  stellar model.
	- $\bullet$  In the absence of rotation, the g-modes arise due to composition gradients.
	- The zero-frequency perturbations are all r-modes.
	- Therefore, as the star rotates, the only inertial modes are the r-modes.
	- The star supports the full set of  $l > |m|$  r-modes, including radial overtones.
- However, neither of these two regimes describe a neutron star.

#### **Globally non-barotropic stars**

 $\bullet$  The *r*-mode frequency is analytic at leading order in  $\epsilon = \Omega/\sqrt{GM/R^3}.$  The solution is obtained by solving the eigenvalue problem for  $\tilde{\omega}_2$ , where

$$
\omega = \omega_0 + \omega_2 + \mathcal{O}(\epsilon^5)
$$

and  $\tilde{\omega}_2 = \omega_2/\sqrt{GM/R^3}$ .

- We explore what happens to the solutions as  $\Gamma_1 \rightarrow \Gamma$ : except for the fundamental  $l = |m|$  solution, all the perturbations diverge and become generic inertial modes.
- Note that their frequencies may no longer satisfy the instability criterion.



**Figure 4:** Eigenfrequencies of the  $l = 2$  r-modes with  $m = 1$  (left panel) and  $m = 2$  (right panel) as  $\Gamma_1 \rightarrow \Gamma = 2$ .

• We develop the r-mode perturbation equations to allow for realistic neutron-star matter with local barotropicity. To do this, we focus on  $l = |m|$  solutions.

**Table 1:** Eigenfrequencies  $\tilde{\omega}_2/\epsilon^3$  of the  $l=|m|$  *r*-modes.

	$m=3$		$m=2$		$m=1$	
k.	BSk19	BSk21	BSk19	BSk21	BSk19	BSk21
	0.4278	0.4282	0.3986	0.3989	0.0000	0.0000
	$-2.6945$	$-0.4610$	$-6.2134$	$-1.2297$	$-18.2599$	$-3.7280$
2	$-5.1463$	$-1.6301$	$-14.6769$	$-3.9823$	$-53.6627$	$-11.9555$
	$-11.5551$	$-3.1305$	$-28.7989$	$-7.6752$	$-94.0447$	$-23.5897$



**Figure 5:** Eigenfunctions of the  $(l, m, k) = (|m|, 3, 1)$  r-modes for BSk19 (left panel) and BSk21 (right panel).



Figure 6: Dimensionless Brunt-Väisälä frequency for BSk19 (left panel) and BSk21 (right panel).

- For fast-spinning neutron stars, we argue that the stratification is weak and the problem resembles the inertial modes, which has been solved in relativity (Idrisy, Owen and Jones [2015\)](#page-12-7). (The next step would be to solve the inertial modes with realistic stratification.)
- For slower spins, the problem will need to be solved numerically from the outset.
- The r-modes with their associated gravitational-wave instability are a promising source for continuous-wave detection.
- However, we are unable to determine the r-modes in general relativity.
- We studied the effects of matter on the spectrum and found that
	- (i) as the star changes from being globally non-barotropic to being barotropic, all but the fundamental  $l = |m|$  r-mode diverge and join the inertial-mode family; and
	- (ii) for realistic nuclear equations of state, the presence of locally barotropic regions results in the same behaviour.
- Because of these effects, we suggest reformulations of the problem to calculate the inertial-mode spectrum with realistic nuclear matter.

### **References**

- <span id="page-12-5"></span><span id="page-12-0"></span>
	- R. Abbott et al. (2021). Astrophys. J. 922.1, 71. DOI: [10.3847/1538-4357/ac0d52](https://doi.org/10.3847/1538-4357/ac0d52).
	- N. Andersson (1998). Astrophys. J. 502.2, 708. DOI: [10.1086/305919](https://doi.org/10.1086/305919).



- N. Andersson and F. Gittins (2023). Astrophys. J. 945.2, 139. DOI: [10.3847/1538-4357/acbc1e](https://doi.org/10.3847/1538-4357/acbc1e).
- <span id="page-12-2"></span>N. Andersson and K. D. Kokkotas (2001). Int. J. Mod. Phys. D 10.4, 381. DOI: [10.1142/S0218271801001062](https://doi.org/10.1142/S0218271801001062).
- <span id="page-12-4"></span>L. Bildsten (1998). Astrophys. J. 501.1, L89. DOI: [10.1086/311440](https://doi.org/10.1086/311440).
- <span id="page-12-1"></span>J. L. Friedman and S. M. Morsink (1998). Astrophys. J. 502.2, 714. DOI: [10.1086/305920](https://doi.org/10.1086/305920).
- F. Gittins and N. Andersson (2023). Mon. Not. R. Astron. Soc. 521.2, 3043. DOI: [10.1093/mnras/stad672](https://doi.org/10.1093/mnras/stad672).
- <span id="page-12-7"></span>A. Idrisy, B. J. Owen and D. I. Jones (2015). Phys. Rev. D 91.2, 024001. DOI: [10.1103/PhysRevD.91.024001](https://doi.org/10.1103/PhysRevD.91.024001).
- <span id="page-12-6"></span>Y. Kojima (1998). Mon. Not. R. Astron. Soc. 293.1, 49. DOI: [10.1046/j.1365-8711.1998.01119.x](https://doi.org/10.1046/j.1365-8711.1998.01119.x).
- <span id="page-12-3"></span>L. Lindblom, B. J. Owen and S. M. Morsink (1998). Phys. Rev. Lett. 80.22, 4843. DOI: [10.1103/PhysRevLett.80.4843](https://doi.org/10.1103/PhysRevLett.80.4843).