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R-modes as probe of Dark matter in neutron stars

(S. Shirke, S. Ghosh et al, Journal of Cosmology and Astroparticle Physics(JCAP) 12 (2023), 008)

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R-modes as source of CGW

 \triangleright Toroidal mode of fluid oscillation in neutron stars for which the restoring force is the Coriolis force.

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- \triangleright Unstable in all rotating stars due to the Chandrasekhar-Friedman -Schutz(CFS) mechanism although dissipation mechanisms can damp and saturate the oscillations. Dissipation timescales determine the instability window.
- \triangleright Spindown of young pulsars, accreting pulsars leading to continuous GW emissions.

R-mode instability window (Chatterjee+ 2006)

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➢ EOS inference from r-modes - **Ghosh, MNRAS 525, 448-454 (2023)**

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- \triangleright GW Searches using LVK data:
- Crab pulsar (Rajbhandari et al. 2021)
- PSR J0537-6910 with $n \approx 7$ (Fesik & Papa 2020, Abbott et al. 2021)
- \triangleright No detection of GW but upper limits on r-mode amplitude were obtained.

Dark matter in Neutron stars

- \triangleright Dark matter ~ 25% of the Universe.
- \triangleright Particle nature of dark matter still unknown.
- ➢ Possibility of Dark matter-admixed neutron star with a DM core or a DM halo.
- \triangleright Primary driving mechanisms:
	- Accretion/Capture
	- Particle Decay
	- DM Seed

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 \triangleright Constraints on the dark matter mass and self-interaction from neutron star observations.

Neutron Decay Anomaly

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- Neutron lifetime:
	- Bottle/Trap Method
	- Beam Method
- Discrepancy in measured lifetime: $\Delta \tau \sim 9.5$ s with 4σ
- Dark decay channel?
- Branching ratio $\sim 1\%$
	- i. e. $\Gamma_{\text{dark}}/\Gamma_{\text{proton}} \approx 1/100$

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PHYSICAL REVIEW LETTERS 120, 191801 (2018)

Dark Matter Interpretation of the Neutron Decay Anomaly

Bartosz Fornal and Benjamín Grinstein Department of Physics, University of California, San Diego, 9500 Gilman Drive, La Jolla, California 92093, USA

EOS model description

● **Hadronic Matter:** Relativistic Mean-Field Model (RMF)

$$
\mathcal{L}_{int} = \sum_{N} \bar{\psi}_{N} \left[\overline{g_{\sigma}} \overline{\rho} - \overline{g_{\omega}} \gamma^{\mu} \omega_{\mu} - \frac{\overline{g_{\rho}}}{2} \gamma^{\mu} \tau \cdot \rho_{\mu} \right] \psi_{N} - \frac{1}{3} b \eta_{N} (g_{\sigma} \sigma)^{3} - \frac{1}{4} c (\overline{g_{\sigma}} \sigma)^{4} \n+ \left[\Lambda_{\omega} (g_{\rho}^{2} \rho^{\mu} \cdot \rho_{\mu}) (g_{\omega}^{2} \omega^{\nu} \omega_{\nu}) + \frac{\zeta}{4!} (g_{\omega}^{2} \omega^{\mu} \omega_{\mu})^{2} \right]
$$

Parameters:

 $\{n_{\rm sat},\,E_{\rm sat},\,K_{\rm sat},\,E_{\rm sym},\,L_{\rm sym},\,m^{\star\prime\prime\prime}\}$

[86] Hornick et al., PRC 98 (2018) [7] Ghosh, D. C., Schaffner-Bielich, EPJA 58 (2022)

● **Dark Matter:** Self-interacting fermionic DM in chemical equilibrium

$$
n \to \chi + \phi
$$
\n
$$
\mathcal{L} \supset -g_V \bar{\chi} \gamma^{\mu} \chi V_{\mu} - \frac{1}{4} V_{\mu\nu} V^{\mu\nu} + \frac{1}{2} m_V^2 V_{\mu} V^{\mu}
$$
\n
$$
\epsilon_{DM} = \frac{1}{\pi^2} \int_{0}^{k_{F_{\chi}}} k^2 \sqrt{k^2 + m_{\chi}^2} dk + \frac{1}{2} G n_{\chi}^2
$$
\n
$$
\left(\text{Vector repulsion}\right)
$$
\n
$$
\mu_{\chi} = \sqrt{k_{F_{\chi}}^2 + m_{\chi}^2} + G n_{\chi}
$$
\n
$$
G = \left(\frac{g_V}{m_V}\right)^2 \qquad n_{\chi} = \frac{k_{F_{\chi}}^3}{3\pi^2}
$$

Cross Section and Shear Viscosity

Cross Section and Shear Viscosity

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Bulk viscosity and the reaction timescale

 \triangleright Bulk viscosity depends on the relaxation timescale of the reaction.

$$
\zeta = P(\gamma_{\infty} - \gamma_0) \frac{\tau}{1 + (\omega \tau)^2}
$$

Experimental constraints : $\Gamma_{\text{dark}}/\Gamma_{\text{proton}} \approx 1/100$

 $Dark$ sector Process:

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- 1. Modified URCA Process :
- 2. Constant timescale :
- 3. Direct URCA process:

$$
n \leftrightarrow \chi + \varphi
$$

$$
n + N \rightarrow p + e^- + \bar{\nu}_e + N
$$

$$
p + e^- + N \rightarrow n + \nu_e + N
$$

 $\tau = 8.88 \times 10^4$ sec

 $n \rightarrow p + e^- + \bar{\nu}_e$ $p + e^- \rightarrow n + \nu_e$

- Two free parameters: G and τ
- **Conserverse Experimental Shirke, Ghosh et al, JCAP 12 (2023) 008**

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radial distance (km)

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R-mode instability window

 \triangleright Boundary obtained by matching the dissipation and GW emission timescales.

$$
\frac{1}{\tau(\Omega_c)}=\frac{1}{\tau_{GW}(\Omega_c,T)}+\frac{1}{\tau_{SV}(\Omega_c,T)}+\frac{1}{\tau_{BV}(\Omega_c,T)}=0
$$

 \triangleright *f-T*—*f - f boundary* \cdot assuming that the power-loss due to the spin-down driven by r-mode instability is equal to the luminosity (both neutrino and photon luminosity) of the star.

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Summary

- ➢ **First investigation of the effect of DM on r-modes of NSs and instability window**
- ➢ **Hadronic model: RMF; DM model: Neutron Decay Anomaly**
- ➢ **Constraints: CEFT, Multi-messenger astrophysics**

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- ➢ **DM SV negligible for T<1010 K**

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- ➢ **Negligible effect of G on r-mode instability window**
- \triangleright **DM** BV significant in cases when $\tau_{DM} = 100 \times \tau_{d-11rca}$ and τ =constant
- ➢ **Pulsar timing data from radio ATNF catalogue and X-ray data of LMXB can be explained in the constant time scale scenario if the timescale is low; else remains incompatible**

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- ➢ **Future: Possible for any other DM model? Microscopic calculation for SV?**